Movement of melt under static and dynamic conditions

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Motivation: Transport of material



Magma chambers and material transport



Cooling of a dyke



To prevent complete crystallization of the dyke new magma has to flow through the dyke at a constant speed (which depends on the dyke diameter and temperature of magma and wall rock).

Minimum velocities: 3.9mm/s – 1m 2.1mm/s – 10m ~7m/h – 14 m/h



Length scales of magmatic transport



Magmatic processes involve a large range of scales

- Initial melt pockets: 10 µm scale (10⁻¹⁵ m³)
- Batholiths 10 km scale (10¹² m³)
- Scale range 10⁹
- Volume concentration factor 10^{27} x

Wetted grain boundaries







Theoretical prediction of melt pocket shapes

Melt penetration



Distribution of melt (Ga) between aluminum crystals and evolution of the thickness of the boundary (synchrotron micro-radiography, Pereiro-Lopez et al., 2003)

Example of grain boundary migration under static conditions



- blue: liquid, green: solid
- ~2% liquid fraction
- wetting angle 10°
- surface energy ss=1
- surface energy sl=0.502

GBM with melt (front-tracking)



Starting grain fabric (white – solid, grey melt) Isotropic surface energies, wetting angle ~10°

Anisotropic surface energies (random orientation), wetting angle 10°

Wetting angles, faceting and permeability

Characterization of the degree of faceting :

$$R = \frac{\Delta \overline{G_{sl}}}{\Delta G_{sl}} = \frac{\sum_{i} \overline{A_{i}} \overline{\gamma_{i}}}{\sum_{i} \overline{A_{i}} \overline{\gamma_{i}} + \sum_{i} \widetilde{A_{i}} \widetilde{\gamma_{i}}}$$

and in case of $\overline{\gamma_i} \approx \widetilde{\gamma_i}$

 $F = \frac{\overline{A}}{\overline{A} - \widetilde{A}}$

 $R = fraction \ of \ free \ energy \ of \ faceted \ surfaces$ $G_{sl} = Gibs \ free \ energy \ at \ solid - liquid \ surface$ $\overline{G_{sl}} = Gibs \ free \ energy \ at \ solid - liquid$ $surface \ of \ faceted \ surfaces$ $\underbrace{\gamma = surface \ energy}_{\overline{A} = total \ surface \ area \ of \ flat \ surfaces$ $\widetilde{A} = tota \ surface \ area \ of \ curved \ surfaces$





(from Yoshino et al., Price et al. etc.)

Melt segregation at high wetting angles (?)

Olivin with FeS melt (static annealing)

(3GPa, 1400°C, 5 vol.% FeS, high wetting angle)



from Walte et al., 2007

Transport of melt with high wetting angle

It apparently works somehow (we can see it in experiments) Must (?) rely on a combination of:

- grain boundary migration (melt behaves passive)
- over- (under-)pressured melt (melt drains into larger melt pockets)



Surface tension and pressure

$$\Delta p = \gamma \left(\frac{1}{R_x} + \frac{1}{R_y}\right)$$

with: $\Delta p = pressure difference(Pa)$ $\gamma = surface tension(N/m)$ $R_{x}(R_{y}) = Radii of curvature(m)$



nano=1e-9 micro=1e-6 mm=1e-3

Pressure difference in a perfect sphere (surface tension = 1 N/m



Radius ~2.5 micron = 2.5e-6 m dp~= 8e5 PaSmall radius generates 3.6x higher pressure difference. This is a driving force!

Radius ~ 9 micron = 9e-6 m dp \sim =2.22e5 Pa

FEM - Underpressured melt pockets

Pressure distribution from different underpressured melt pockets



Getting dynamic: pressure distribution around a single melt pocket



FEM – pressure distribution

Pressure distribution around two closely spaced melt pockets (with and without underpressured melt)





•Pressure concentration leads to:

- pressure solution
- development of cracks



Pressure solution



 Solubility (C) of most minerals is pressure-dependent:

 $C_{equilibrium} = C_0 + a \cdot P$

- Material dissolves at high-P points
 - Grain boundaries normal to $\sigma_{\scriptscriptstyle max}$
- Material precipitates at low-P points
 - Grain boundaries normal to $\sigma_{\mbox{\tiny min}}$

Material diffuses from high pressure to low pressure

change of wetting angle
change of melt pocket shape -> change of pressure concentration

Development of cracks

Cracks will develop if:

• the energy release due to cracking is higher than the energy needed to make 2 new surfaces.

Energy increase due to 2 new surfaces of mineral 1 and 2

$$\gamma_1 l + \gamma_2 l \leqslant \frac{\pi \sigma^2 a t}{\epsilon}$$

Strain energy release from cracking (Griffith theory on (brittle) cracking)

• the orientation of the melt pocket is preferable for cracks









Conclusion for melt transport

- <u>Under static conditions melt moves</u>
 - wetting angle
 - by GBM (passive movement)
 - by pressure solution
 - cracking (from pressured melt pockets)
- Under dynamic conditions melt also moves by
 - pressure concentrations at melt pocket tips σ_1 enhances crack formation (will be parallel to γ_1)
 - The orientation of the melt pockets with respect to the pressure (or better applied forces) is very important and plays a major role during the transport of melt!

Outlook

- Couple solid and fluid mechanics in a numerical simulation
- Couple different types of numerical models
 (FEM, Finite Difference, Front tracking)
- Simulate a "real" microstructure (lots of different melt pockets)

Behavior of melt during grain boundary migration and simple shear



Viscosity contrast 1:100

Developing networks





Never fully connected network: transient local connectivity

Local transient flow events at all scales