

29/06/2010

Soutenance de thèse

Comportement hydro-thermique d'un écoulement de fluide dans une fracture rugueuse :

Modélisation et application à des massifs fracturés

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Directeur de thèse : Jean Schmittbuhl

Co-encadrant : Renaud Toussaint

Thèse présentée en vue d'obtenir le grade de :
Docteur de l'Université de Strasbourg

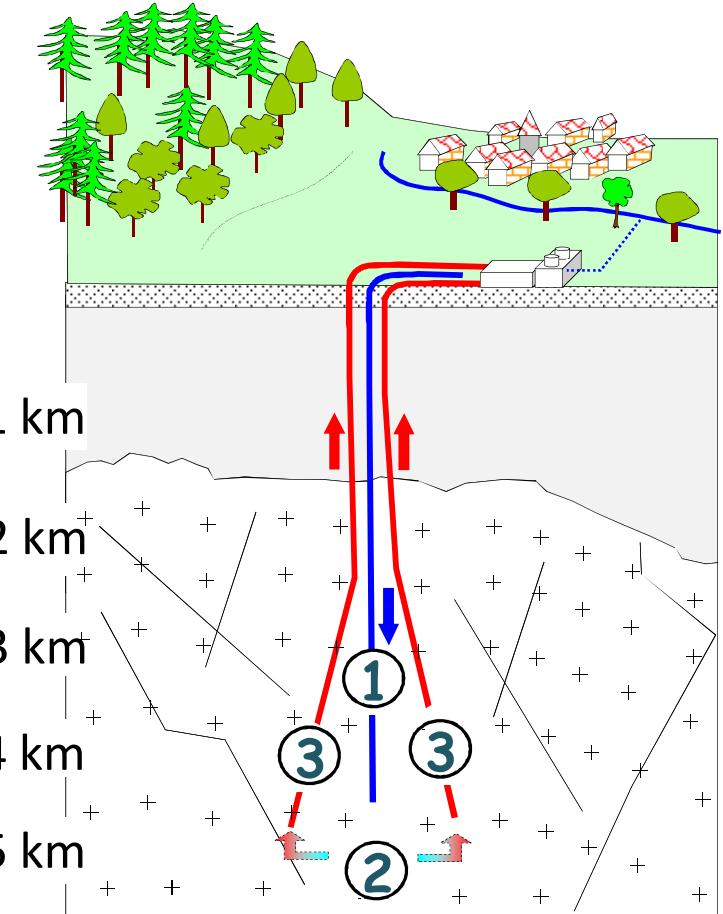
Discipline : Sciences de la Terre et de l'Univers
Spécialité : Géophysique



Geothermal background

- Thermal exchanges between a **hot** fractured rock and a **cold** fluid

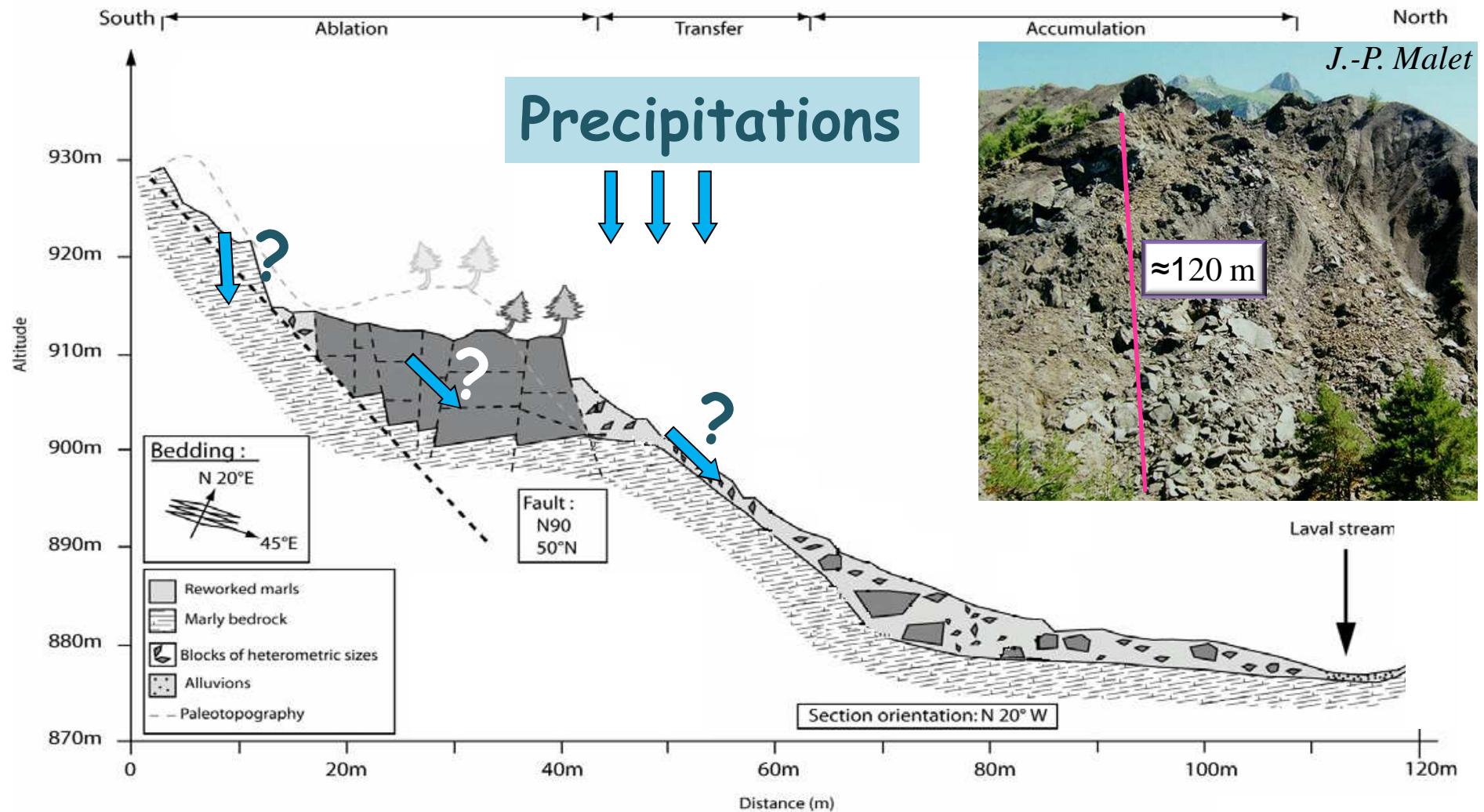
- Deep geothermal systems
 - > “Enhanced Geothermal Systems”
 - Soultz-sous-Forêts (Alsace, France)
 - Cooper Basin (Australia)
 - > Example of parameters
 - Hydraulic flow : 25 l/s
 - Temperature at injection : 60° C
 - Temperature at pumping : 180° C



A. Gallien, d'après documents AREVA

Landslide background

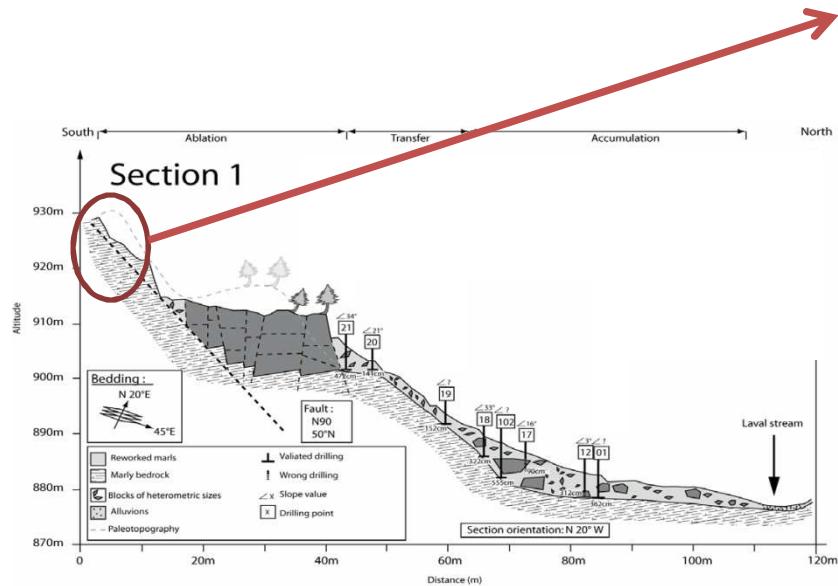
○ Influence of water on landslide triggering



Draix, (Alpes de Haute-Provence, France)

M. Fressard, 2009

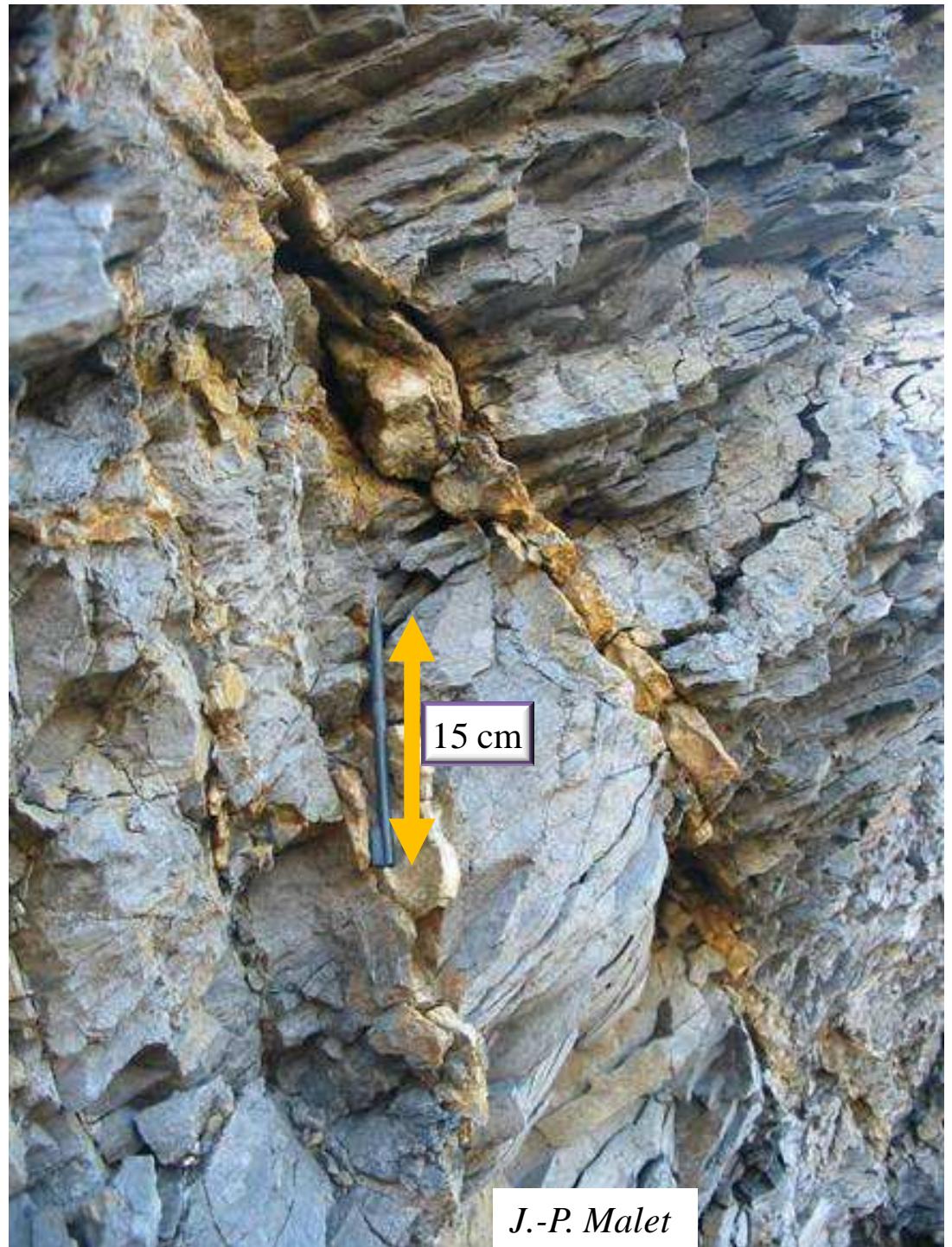
Landslide background



Altered rock

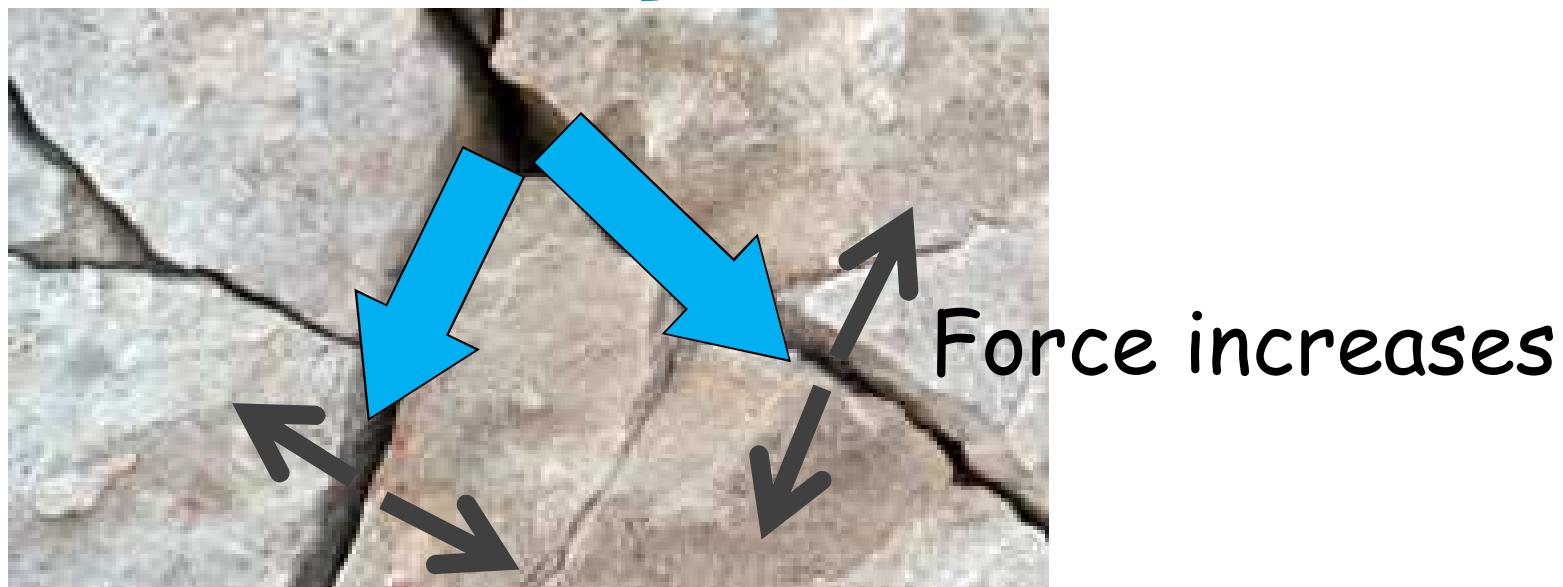
→ Fluid inside fractures

Permeability of
Draix bedrock?



Influence of water on landslides triggering (Not exhaustive)

- Gravity: material full of water is heavier
 - Pore pressure
 - Chemical processes
 - > Rheology/dissolution
 - > Sealing of fractures
- } {
- Less friction
 - More fractures
 - Larger fractures

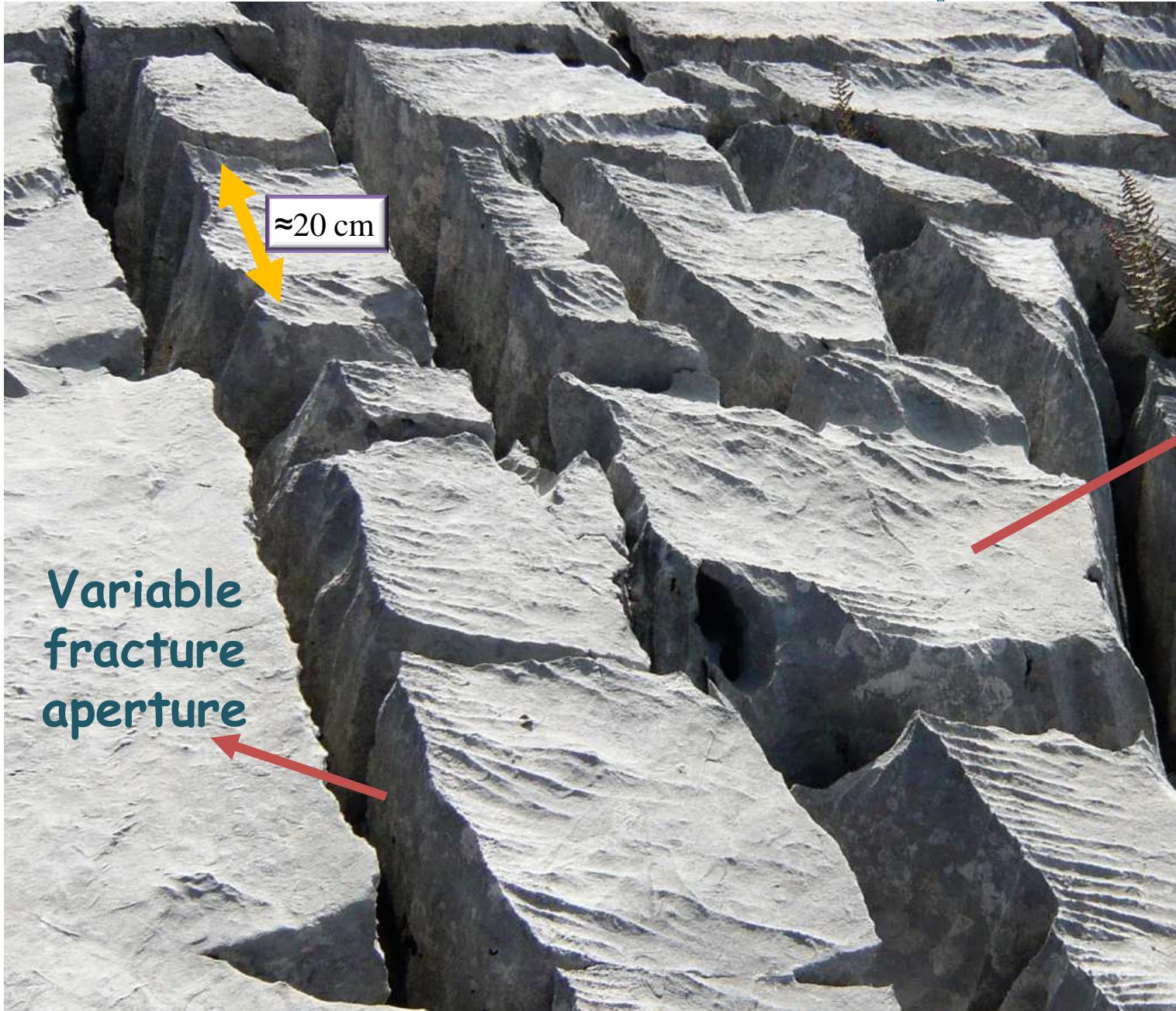


Fracture network



*D.D. Pollard
Limestone bed, England.*

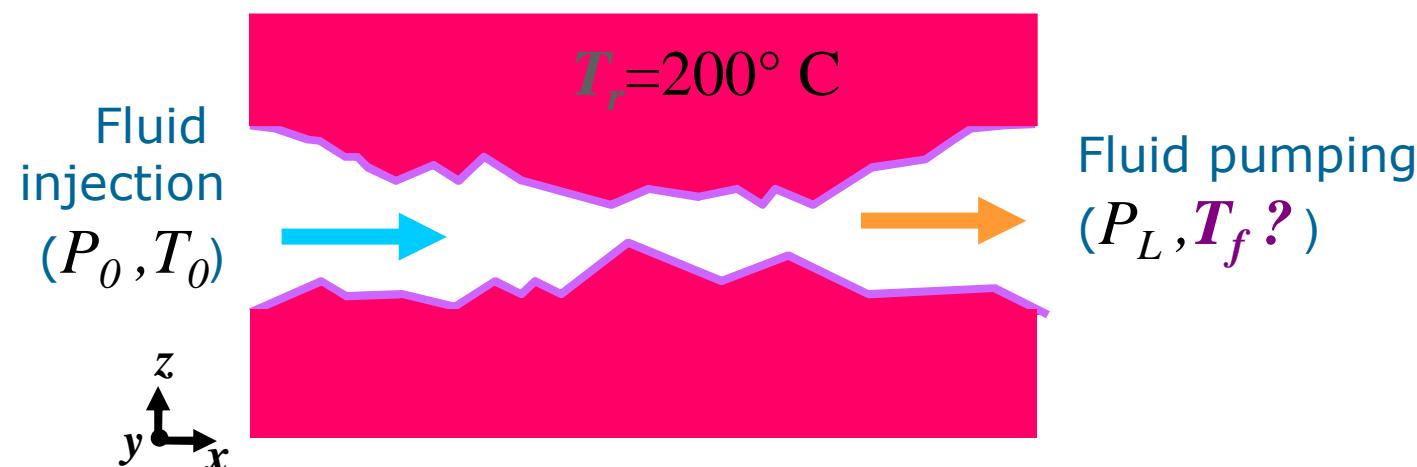
Variable fracture morphology



S. Näff,
Swiss
Speleological
Society.

Questions

- Morphology of fractures?
- Effect of the morphology of fractures on the
 - > Hydraulic flow?
 - > Heat exchange between fluid and rock?



Scale: individual fracture

Outline

○ Morphology

- > Measured on natural fractures (Draix borehole cores)
- > Characterization
- > Synthetic fractures

○ Hydraulic models

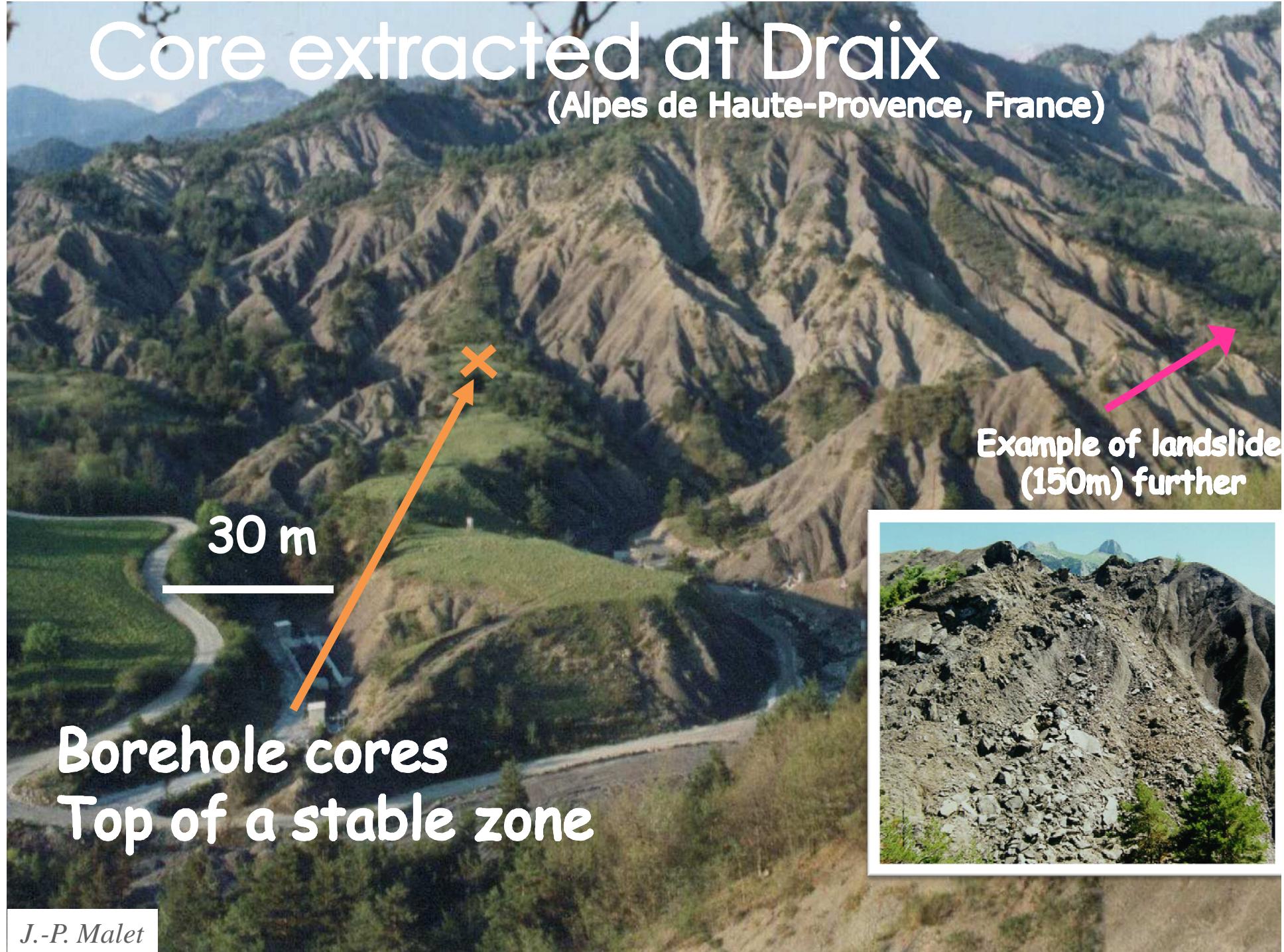
- > Finite differences (methode 1)
 - Hypotheses
 - Results
 - Application to Draix
 - Limits
- > Lattice Boltzmann (methode 2)
 - Bases
 - Implementation

○ Hydro-thermal models

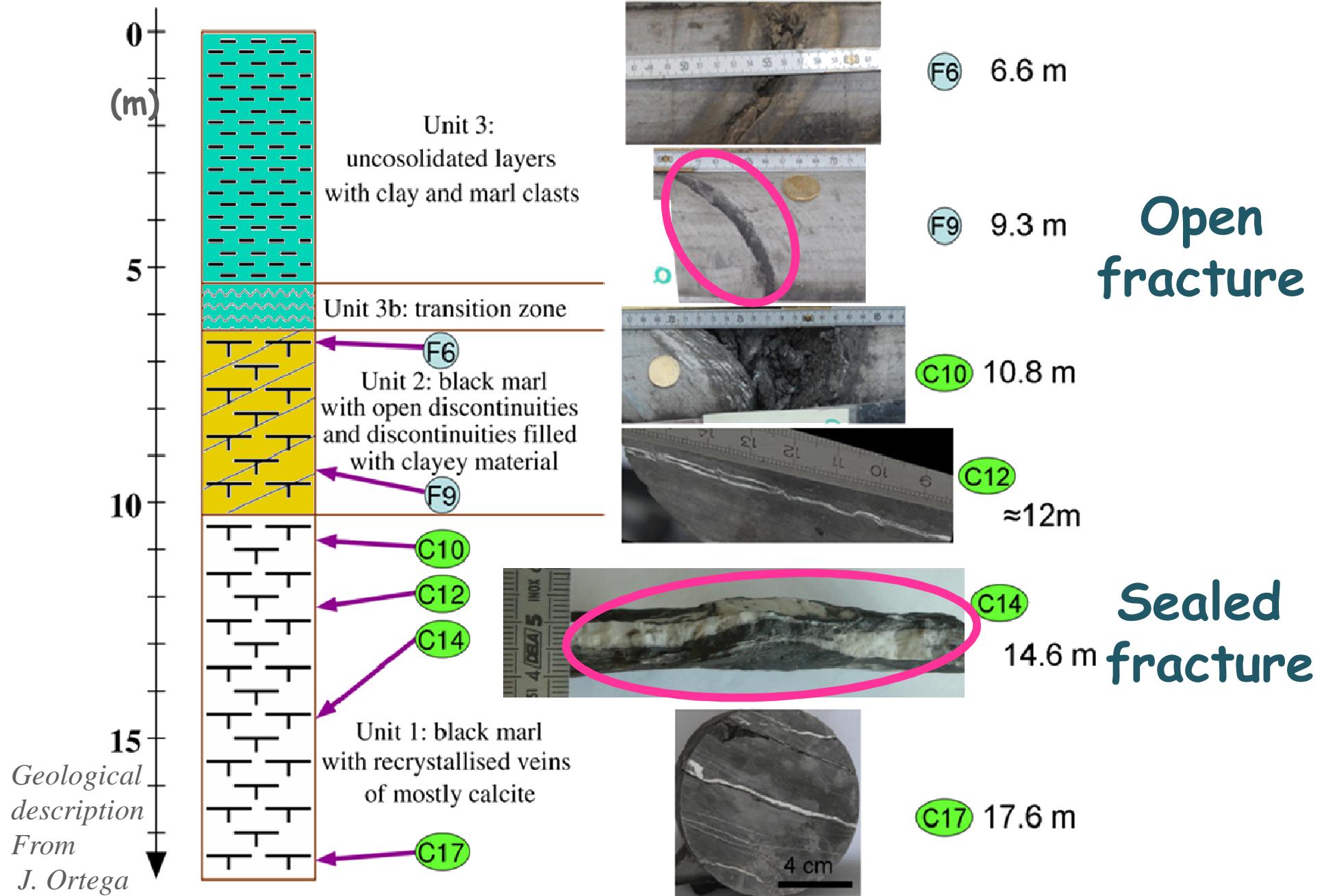
- > Finite differences model
 - Hypotheses
 - Results
 - Limits
- > Lattice Boltzmann (LB) model
 - Bases
 - Implementation

Core extracted at Draix

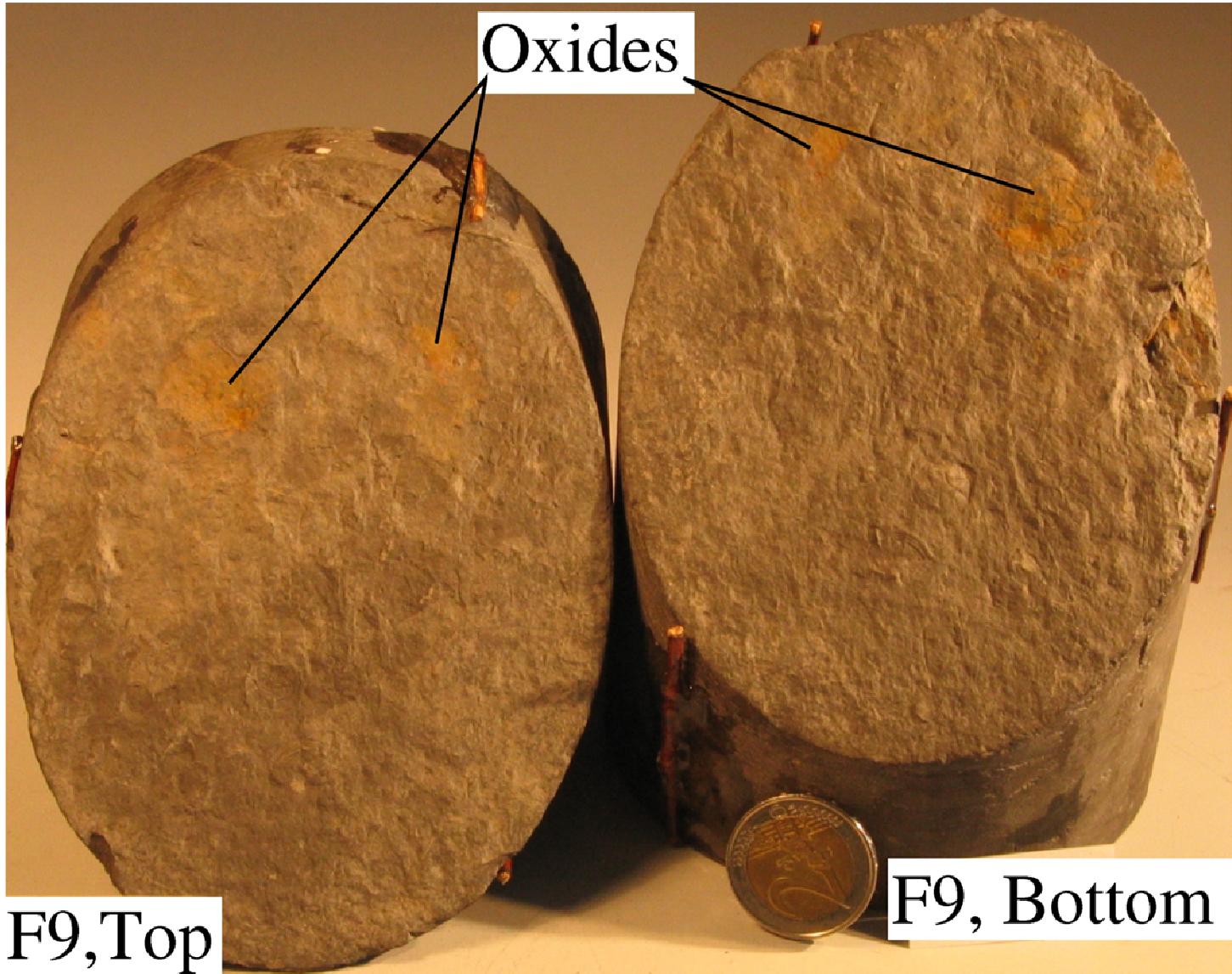
(Alpes de Haute-Provence, France)



Data: Drilled core

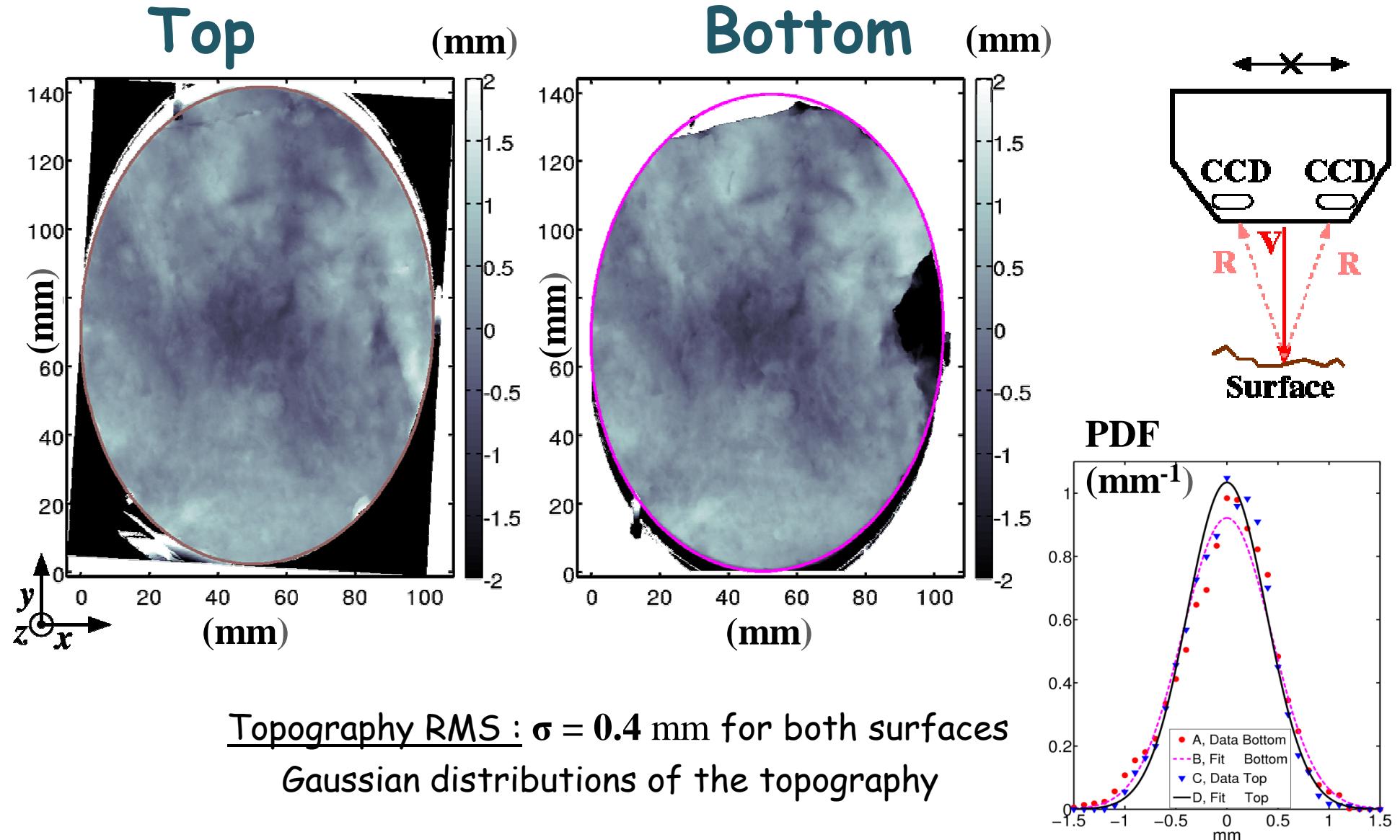


Open fracture



Roughness of the topography

- Optical profiler (vertical precision $\sim 1\mu\text{m}$)

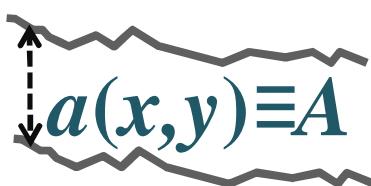


How opened are the fractures?

- Contact ? Mean aperture A ?
- Variability of the aperture ?
- Scale properties of the surfaces
 - > Independence of both surfaces?

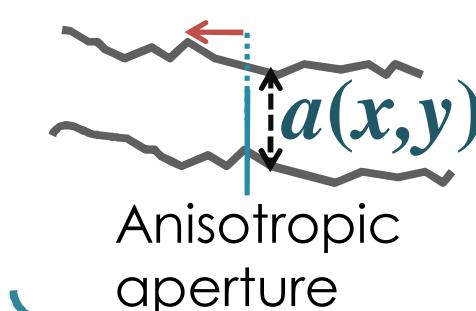
Correlated surfaces

Identical surfaces
Mode I



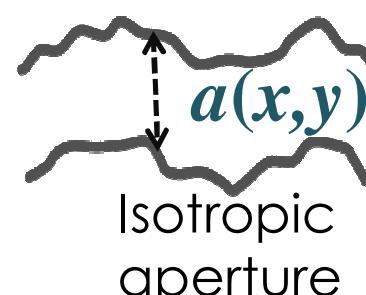
A : constant

Identical surfaces
Mode I+II



Anisotropic aperture

~ Identical surfaces
Small scale noise
Mode I



Isotropic aperture

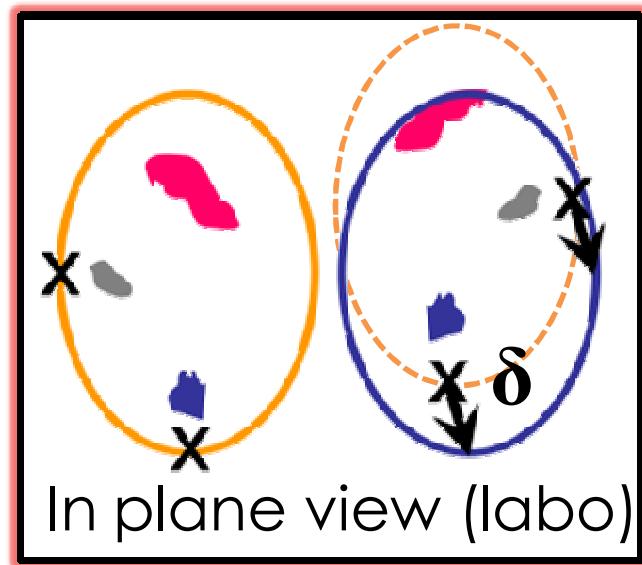
- > Correlated surfaces at large scales
- > Independent surfaces at small scales

Independent surfaces

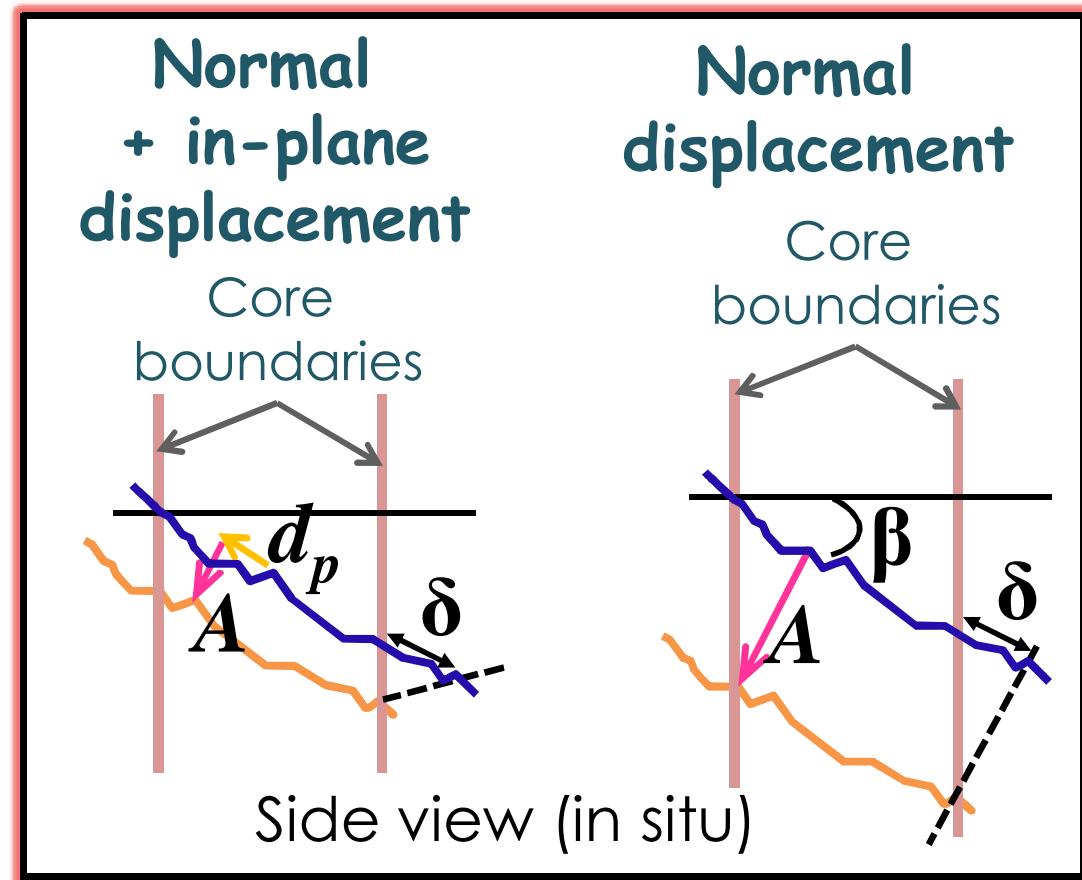


Reconstruction of aperture: open discontinuity

- Similarities of the sides → δ obtained
- Assumptions about the displacement:



In plane view (labo)

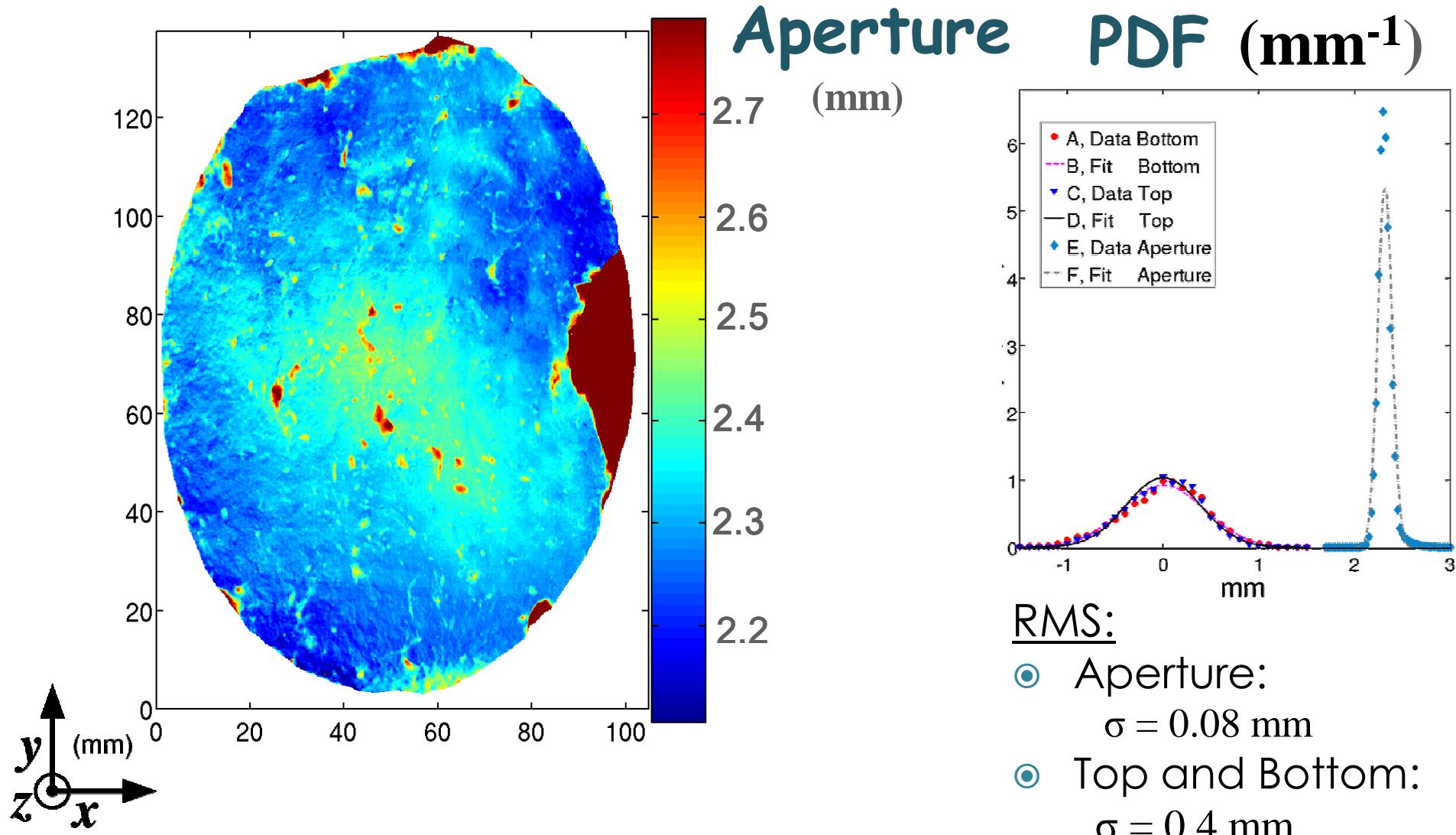


$$\delta = A \cdot \tan(\beta) + d_p$$

- > Pure in-plane displacement → $A = 0$; $d_p = \delta$
- > Pure normal displacement → $A = \delta / \tan(\beta) \approx 2.3 \text{ mm}$; $d_p = 0$

Reconstruction of aperture: open discontinuity

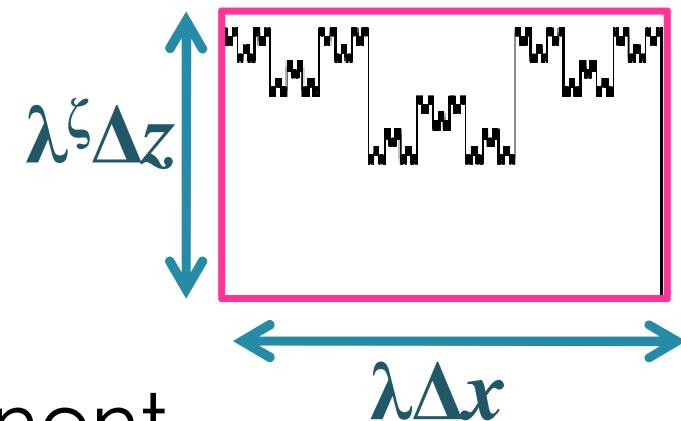
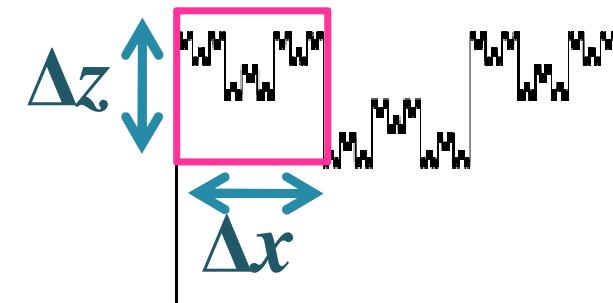
Hypothesis: normal displacement



Self affine topography

- Spatial property:

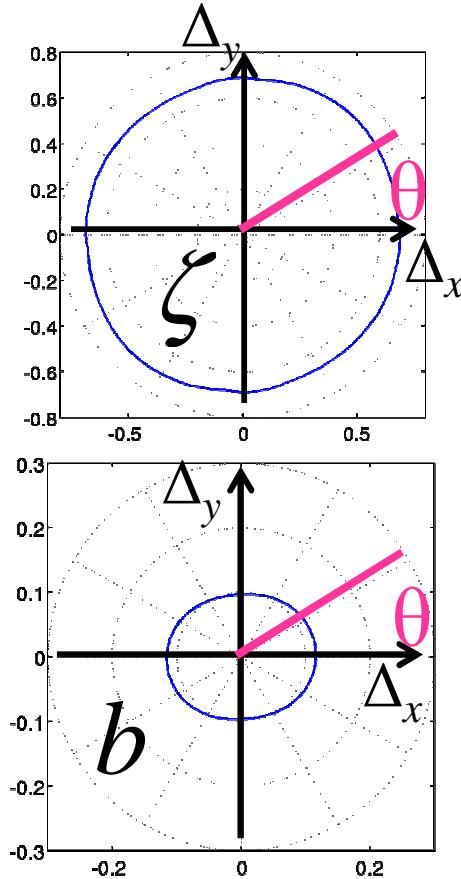
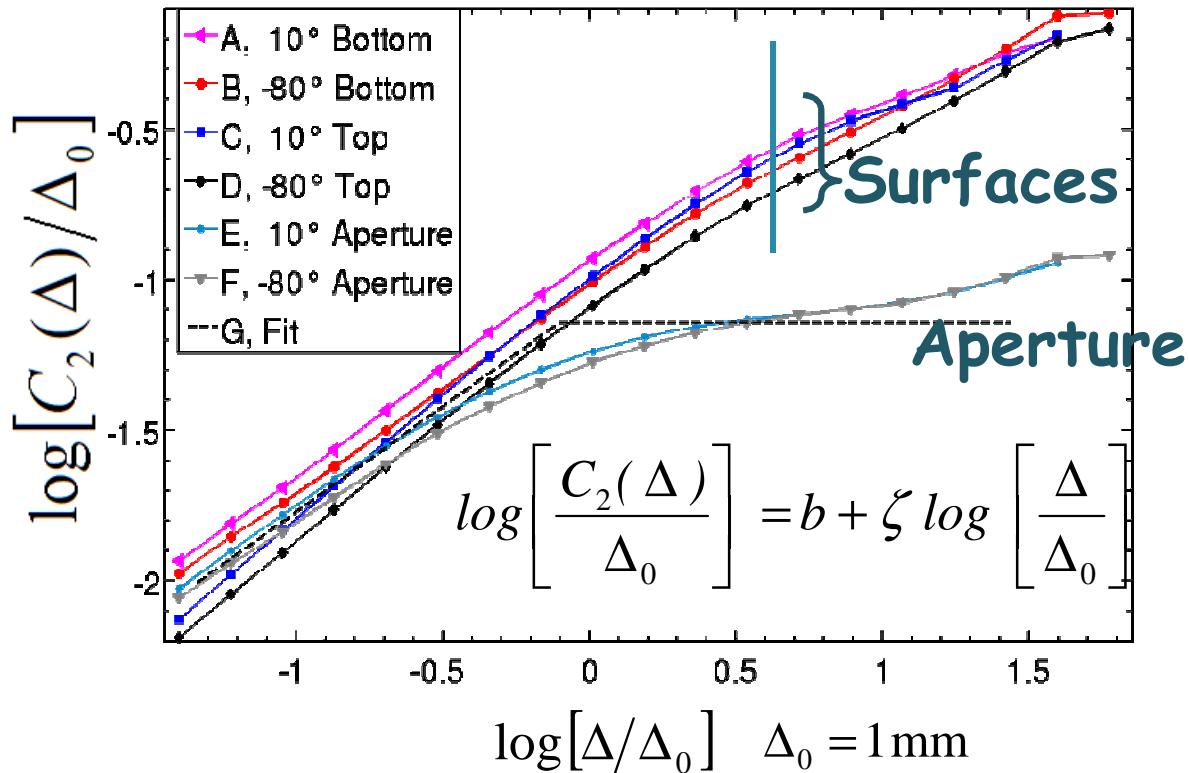
- > Statistical spatial correlation
- > Anisotropic fractal
- > Statistically invariant under:
 - $\Delta x \rightarrow \lambda \Delta x$
 - $\Delta y \rightarrow \lambda \Delta y$
 - $\Delta z \rightarrow \lambda^\zeta \Delta z$ (*for any* λ)



- Roughness (or Hurst) exponent

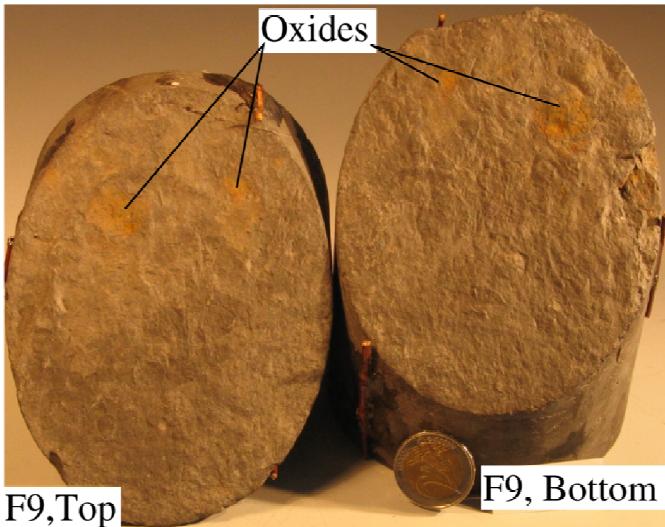
$$\zeta \approx 0.7 - 0.8$$

Autocorrelation C_2 of the topography/aperture

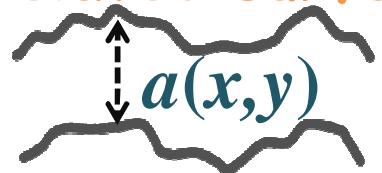


- Bottom and top topographies
 - > Self affine from 0.06 to 7 mm
 - > $\zeta \approx 0.70$ (bottom)
 - > $\zeta \approx 0.75$ (top)
- Aperture
 - > More or less self-affine if $\Delta < 1$ mm
 - > $\zeta \approx 0.6 - 0.7$
 - > More or less Uncorrelated if $\Delta > 1$ mm

Open fracture Aperture measurement

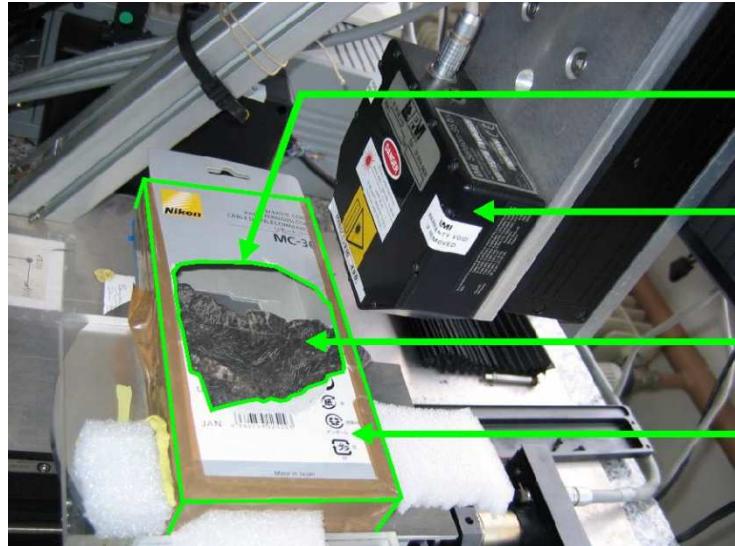


Correlated surfaces

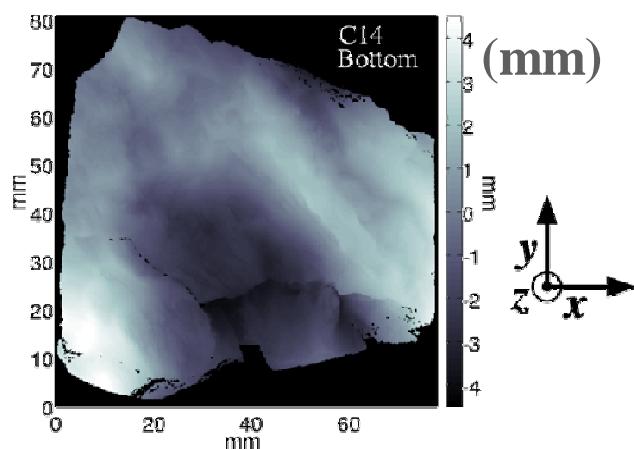
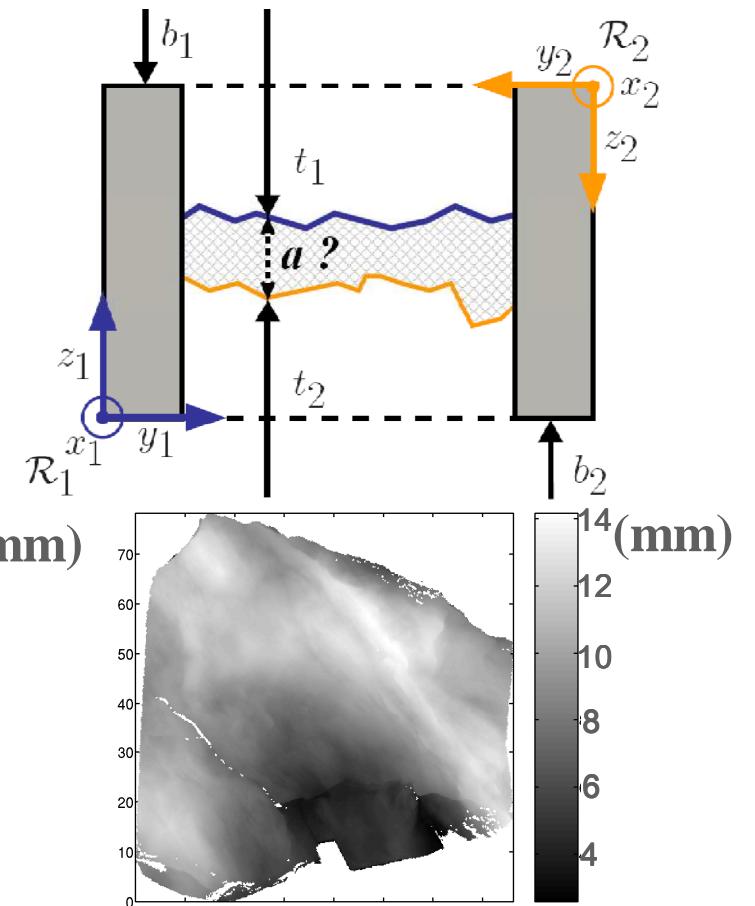


- > No anisotropy of
 - the surfaces
 - the aperture
- > Correlated surfaces at large scales
- > Independent surfaces at small scales
- > Self affine model of the aperture at small scales

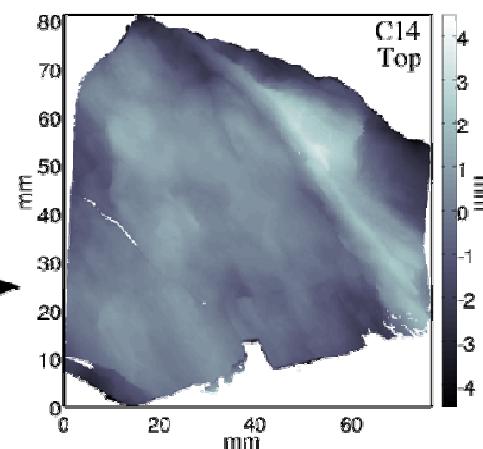
Aperture reconstruction Sealed discontinuity



Window
Profiler
C14 sample
Box

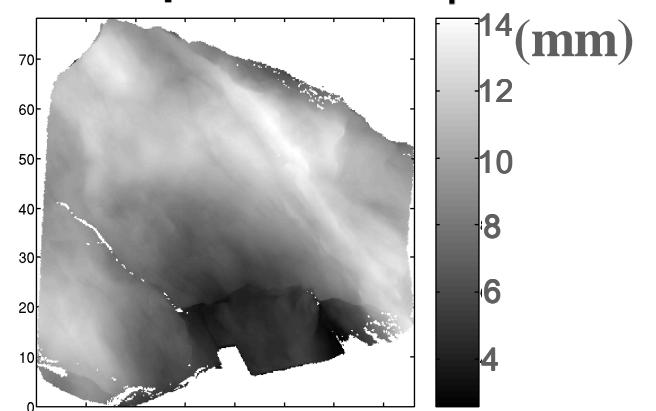


Bottom

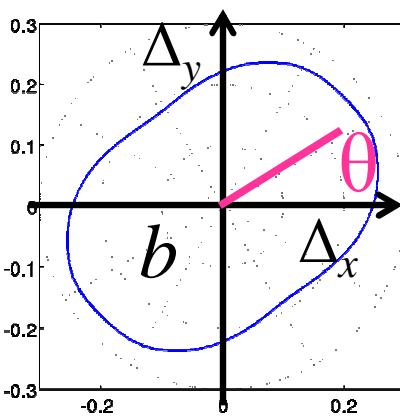
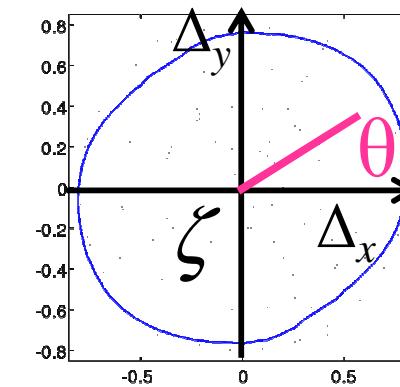
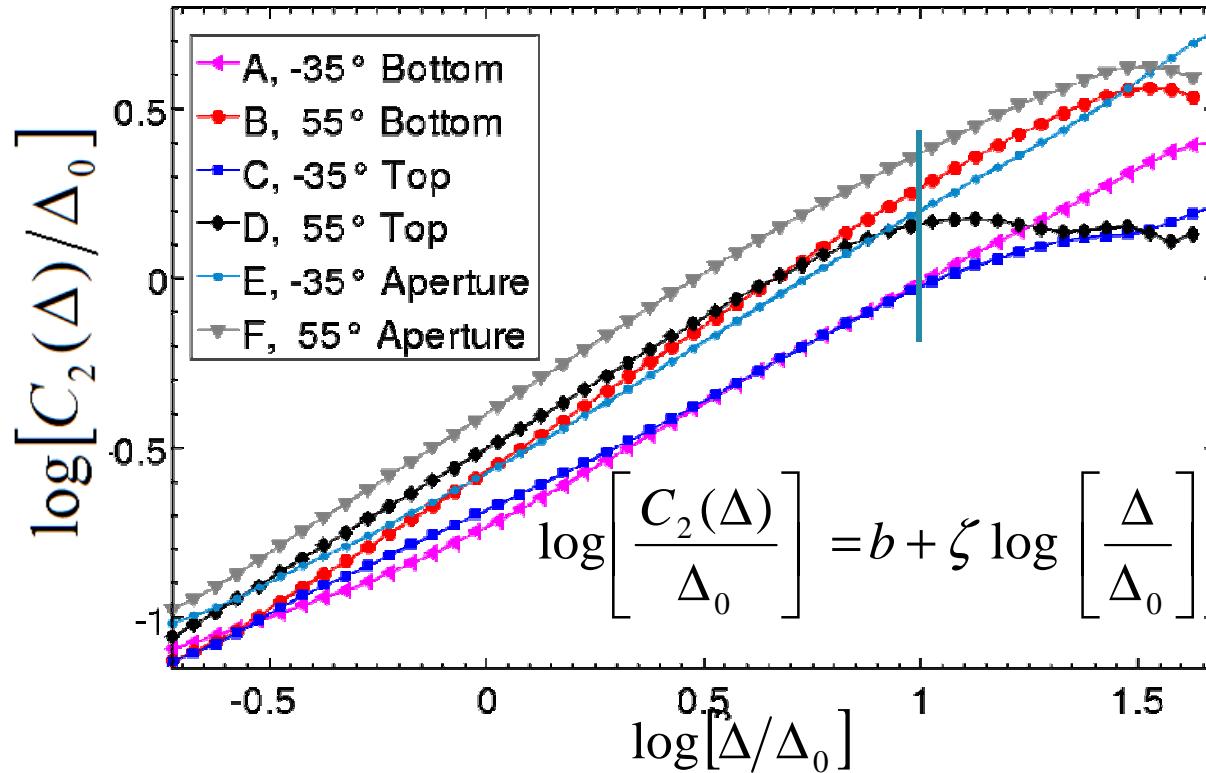


Top

Aperture $\sigma / A > 0.45$



Autocorrelation C_2 - Sealed fracture



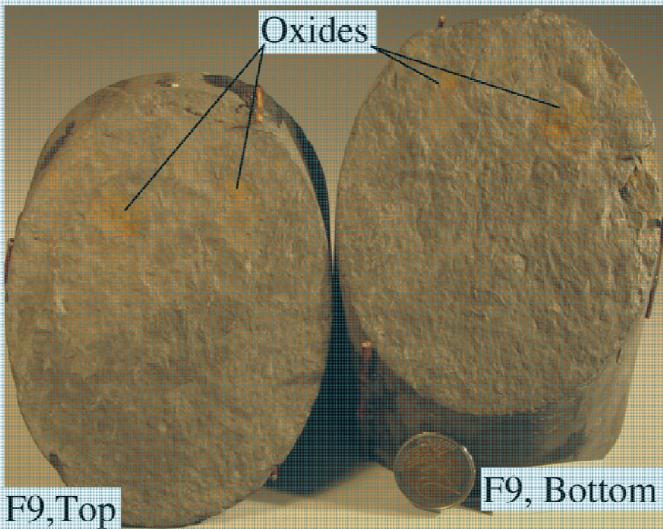
$$\Delta_0 = 1 \text{ mm}$$

Bottom, top and aperture:

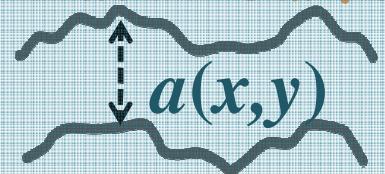
- > Self affine from 0.3 to 10 mm
- > Anisotropy of b and ζ (visible despite errors bars)
 - $\zeta \approx 0.65 - 0.8$ (top)
 - $\zeta \approx 0.7 - 0.85$ (bottom)
 - $\zeta \approx 0.7 - 0.85$ (Aperture)

Aperture measurements

○ Open fracture



Correlated surfaces



- > No anisotropy of the surfaces
- > No anisotropy of the aperture
- > Correlated surfaces at large scales
- > Independent surfaces at small scales

○ Sealed fracture



Independent surfaces

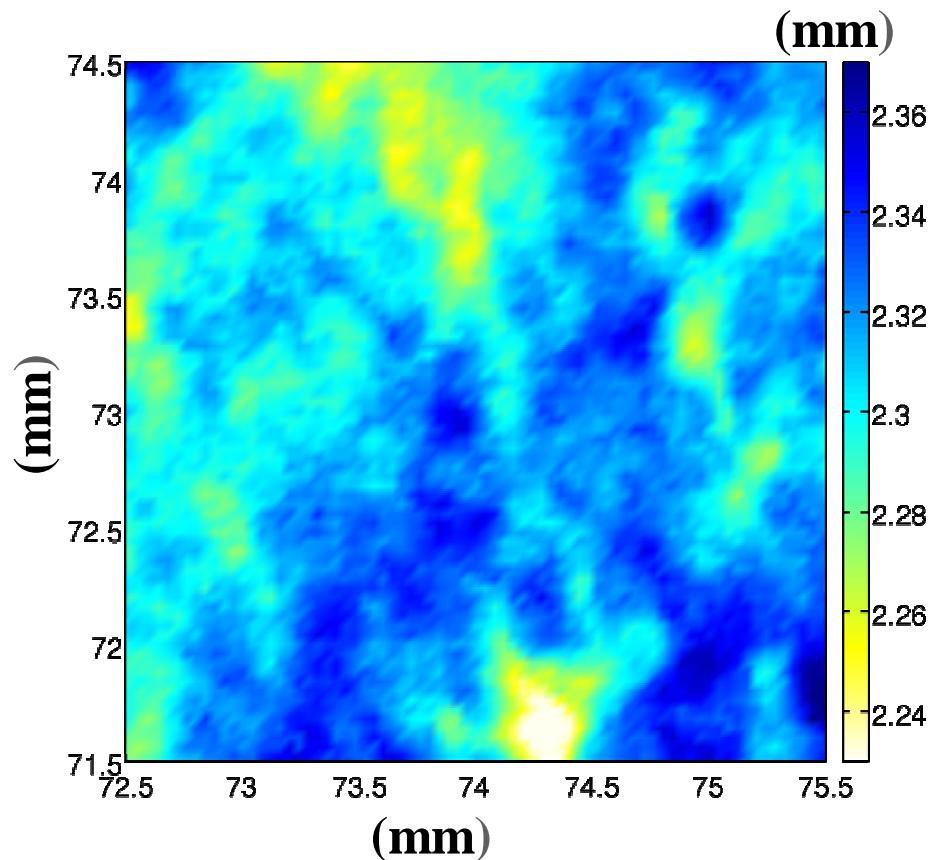


- > Anisotropy of the surfaces
- > Anisotropy of the aperture
- > Independent surfaces

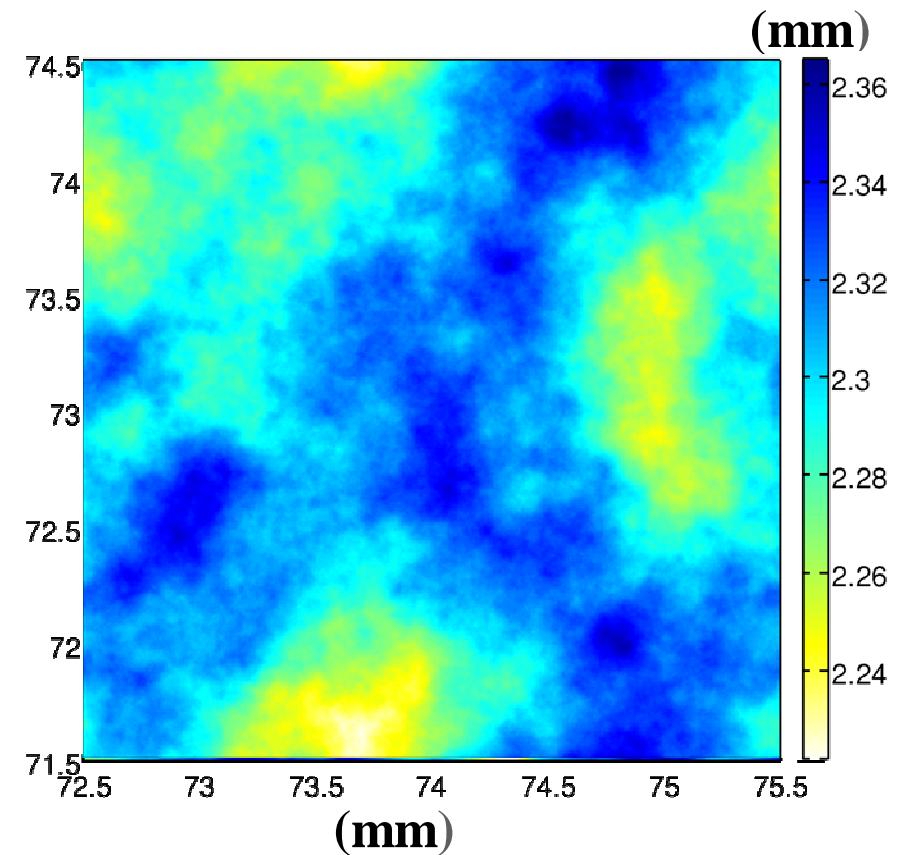
Both are self-affine at small scales

Aperture model

○ Natural aperture



○ Self-affine aperture



$$\zeta = 0.8$$

Hydraulic flow - meth.1: finite differences

- $\rho \left(\frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \vec{\nabla} \vec{v} \right) = -\vec{\nabla} p + \eta \Delta \vec{v}$

*Permanent
Laminar*

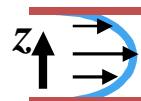
- Stokes $\vec{\nabla} p = \eta \Delta \vec{v}$

Hydraulic lubrication

$$v_z(x, y, z) = 0$$

- Parabolic profile:

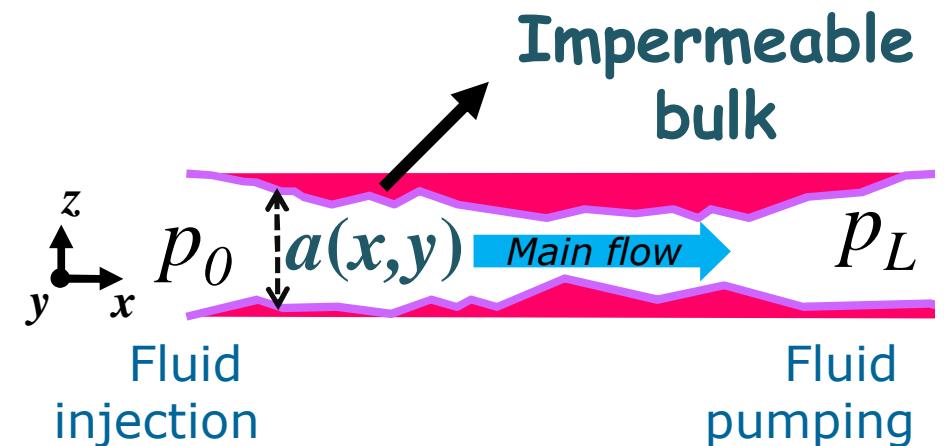
$$\vec{v}(x, y, z) = \frac{\vec{\nabla} p}{2} \cdot (z - z_1) \cdot (z - z_2)$$



Incompressibility

- Reynolds equation:

$$\vec{\nabla} \cdot (a(x, y)^3 \vec{\nabla} p) = 0$$



Notations:

Fluid density: ρ

Dynamic viscosity: η

Velocity: $\vec{v}(x, y, z)$

Pressure: $p(x, y, z)$

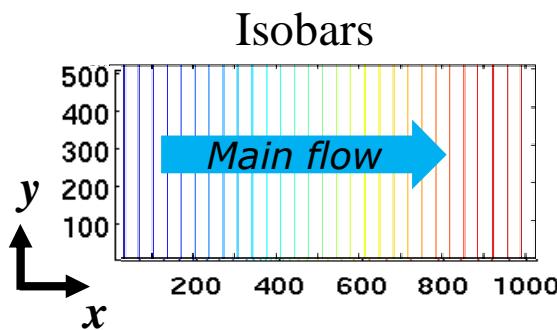
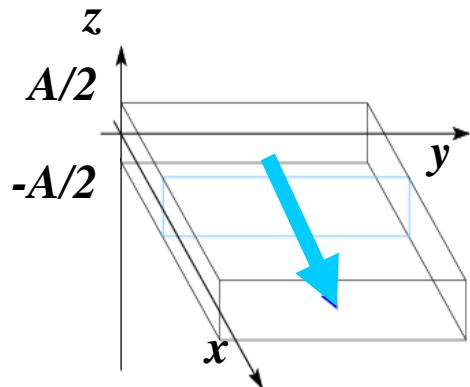
in 2D

Solving: Finite differences + biconjugate gradient method

Hydraulic aperture H

- Parallel plates

➤ *Analytic solution*

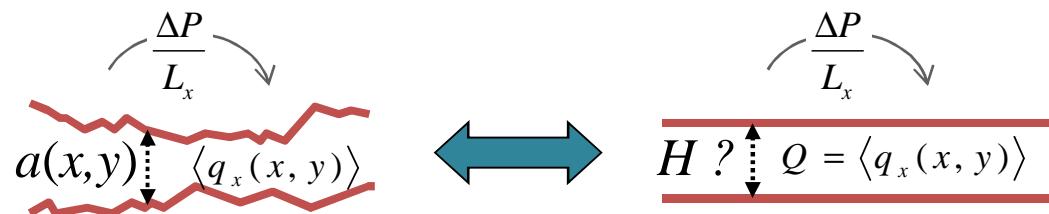


Hydraulic flow

$$\vec{q}_{\parallel}(x, y) = -\frac{\Delta P}{L_x} \frac{A^3}{12\eta} \hat{x}$$

$$\vec{q}(x, y) = \int_a \vec{v}(x, y, z) dz = -\frac{a^3}{12\eta} \vec{\nabla}_2 p$$

- Variable aperture



Geometrical aperture:
 $A = \langle a(x, y) \rangle$

$$Q = -\frac{\Delta P}{L_x} \frac{H^3}{12\eta}$$

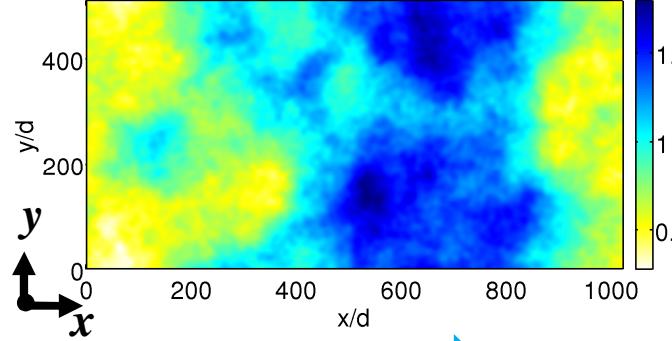
Hydraulic aperture:

$$H = \left[-Q 12 \eta \frac{l_x}{\Delta P} \right]^{1/3}$$

$$H \neq A$$

$$\langle a^{-3} \rangle^{(-1/3)} < H < \langle a^3 \rangle^{(1/3)}$$

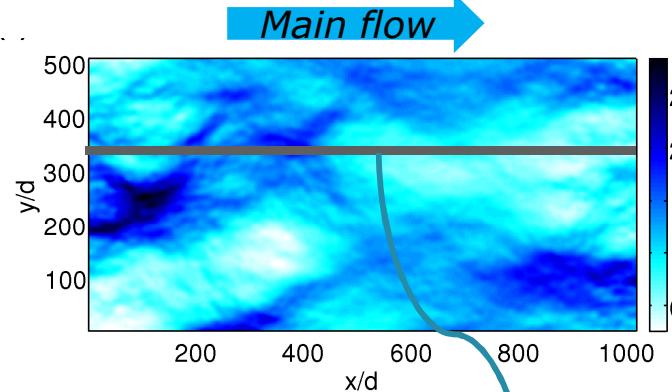
Illustration...hydraulic result



Rough apertures
 $a^*(x, y) = \frac{a(x, y)}{A}$

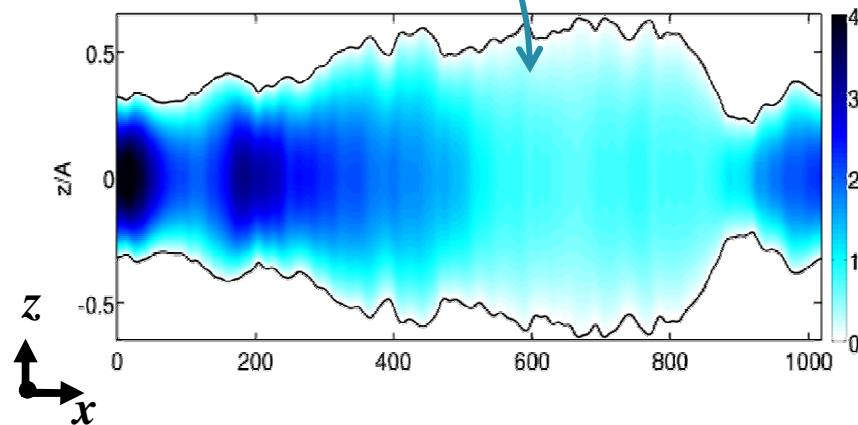
Example

A	1 mm
σ	0.35 mm
$L_x \times L_y$	1 x 0.5 m ²



2D-flow norm

$$\|\vec{q}^*(x, y)\| = \frac{12\eta L_x \|\vec{q}(x, y)\|}{\Delta P A^3}$$

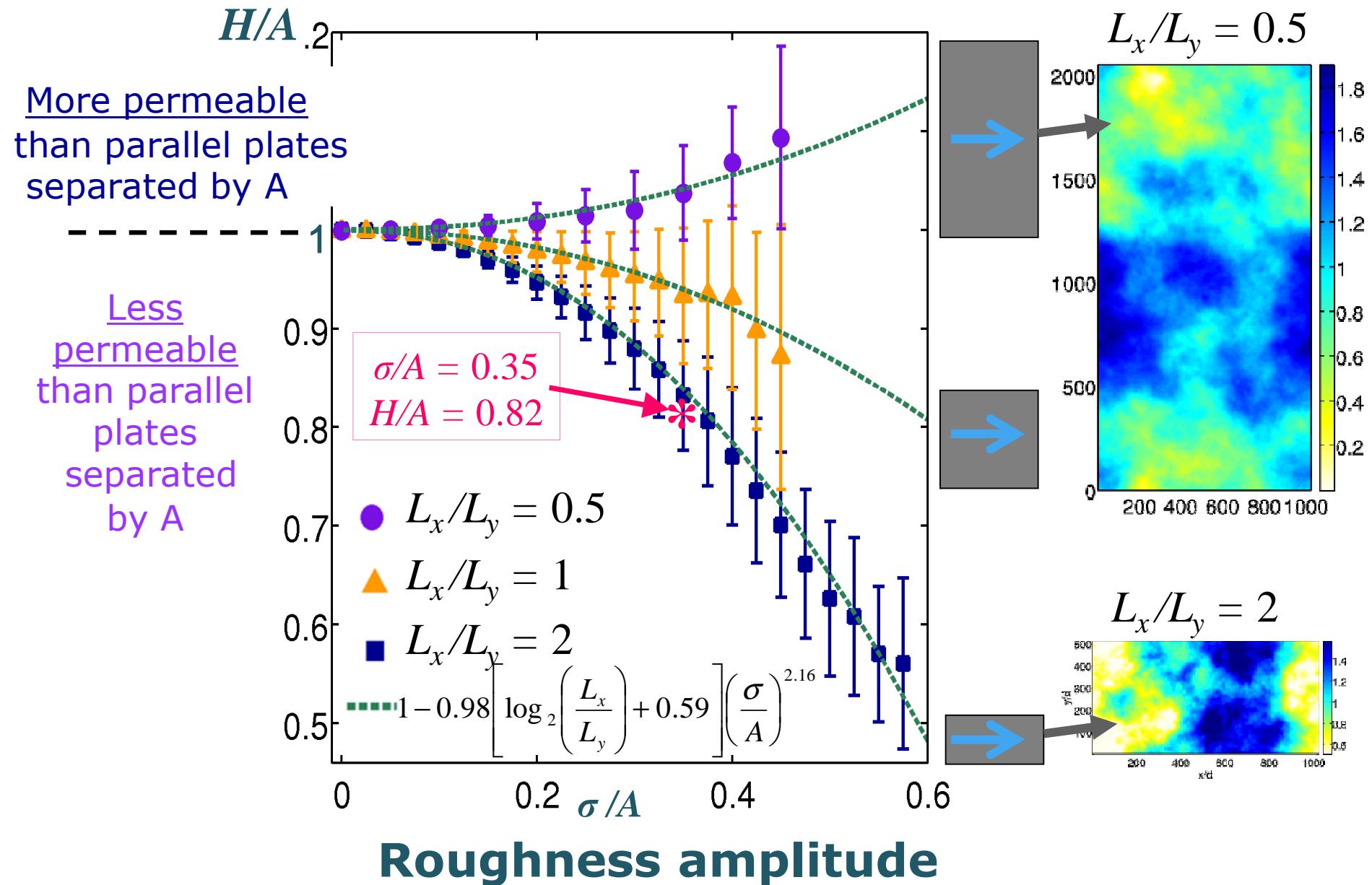


Velocity
 $v^*(y/d=302)$

$$\text{Re} = \frac{\text{inertial forces}}{\text{viscous forces}} = 0.23$$

Dyn. visc. η (10 bar; 100°C)	3.10^{-4} Pa.s
$\partial p / \partial x$	250 Pa/m
H	0.82 mm
$q^* = 1$	$q = 7 \cdot 10^{-5}$ m ² /s
$v^* = 1$	$v = 7 \cdot 10^{-2}$ m/s

Statistical results: hydraulic apertures H



Permeability of the Draix fractures

- Like the open aperture:

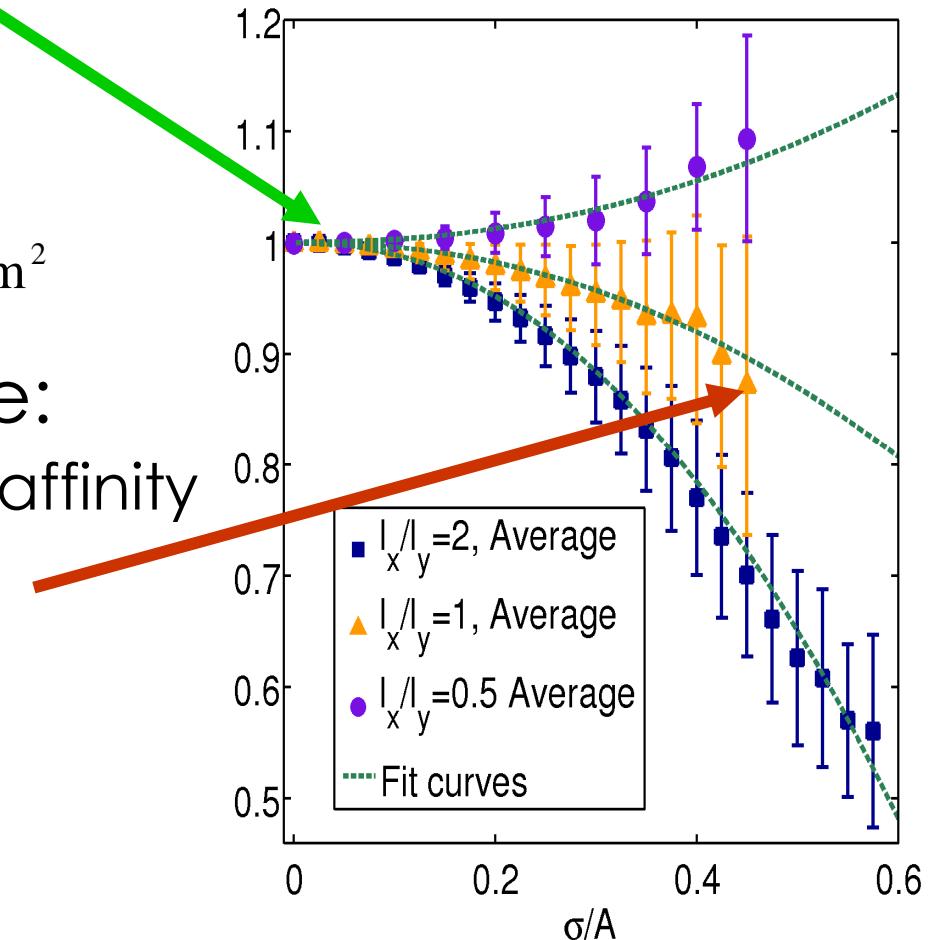
- > About flat aperture : $\sigma /A < 0.04$
- > $H = A = 2.3 \text{ mm}$
- > Permeability:

$$k = \frac{H^3}{12} = 1.10^{-9} \text{ m}^2$$

- Like the sealed aperture:

- > At observed scales : self-affinity
 - $\sigma /A > 0.45$
 - $A \approx 1 \text{ cm}$
 - $H'/A \approx 0.9 \Rightarrow H' \approx 0.9 \text{ cm}$
- > Permeability:

$$k' = \frac{H'^3}{12} = 6.10^{-8} \text{ m}^2$$



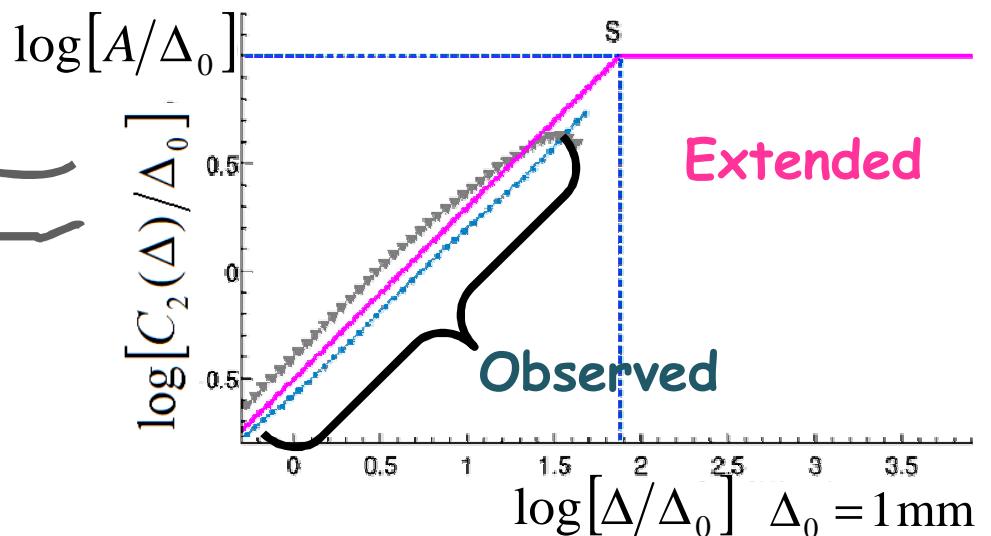
Permeability of Draix bedrock

How to extend fracture models at large scales ?

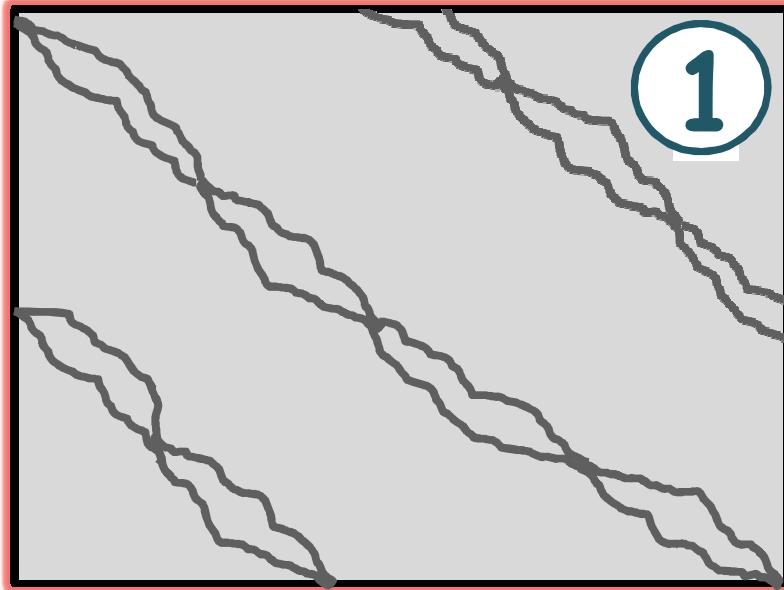
- Linear density fractures
 - > Open fracture: 1 per m (core observation)
 - > Sealed fracture which will reopen
 - Chemical conditions ??
- Extension of the scaling law
 - > What we know:
 - Self-affinity can't be valid at very large scales



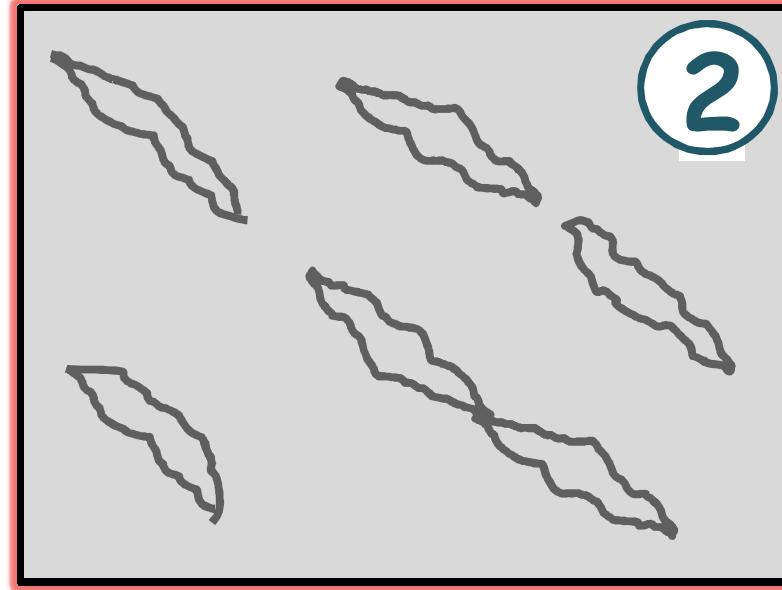
Interpenetration
Not physical !



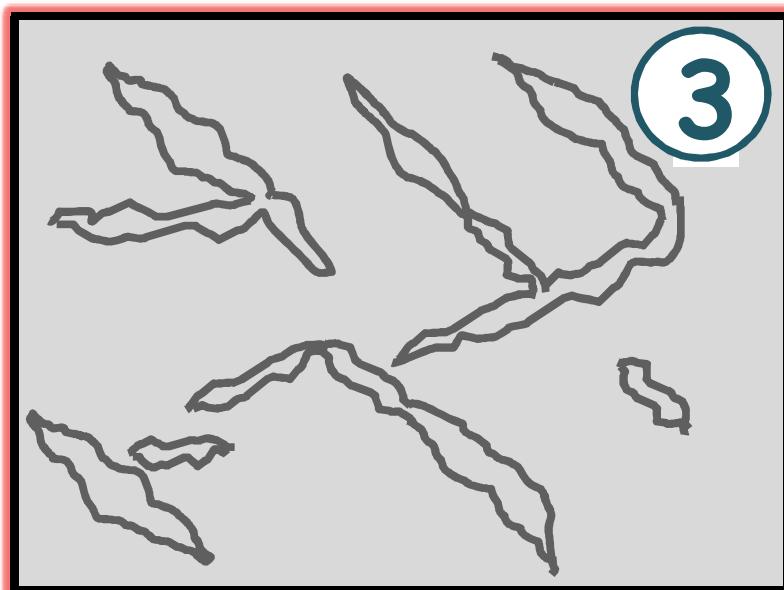
Network connectivity



1



2



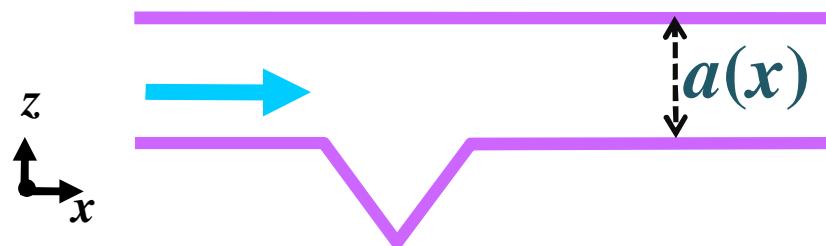
3

○ Are the observed fractures representative of the large scale permeability ?

- 1 Connected network → yes
Local contact
- 2 Disconnected network
Bulk permeability ?
- 3 Complicated network ! → No

Off lubrication regime ?

- Effect of sharp morphology ?
 - > Effect of contact zones ?
 - > Corner



- Higher velocities
- Full Navier Stokes equation

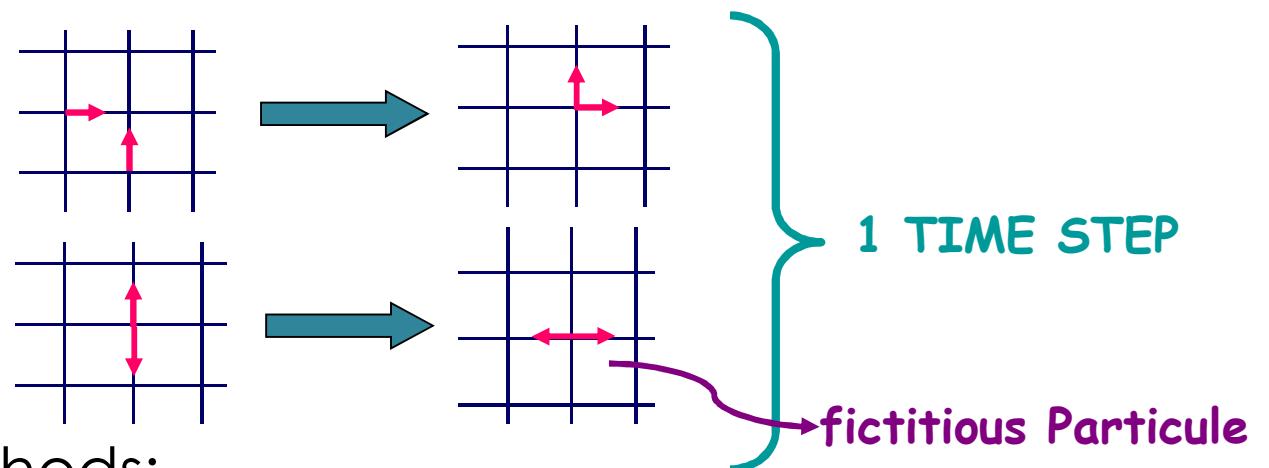
$$\rho \left(\frac{\partial \vec{v}}{\partial t} + \vec{v} \cdot \vec{\nabla} \vec{v} \right) = -\vec{\nabla} p + \eta \Delta \vec{v}$$

- Lattice Boltzmann (LB) methods
- Transient regime
- Equation to be solved in 3D

Some principles of the LB methods

- Comes from Lattice gas methods:
 - > Space discretized with a lattice
 - > Discrete time, discrete velocity directions
 - > Fictitious particles, 1 particle/node in a given direction
 - Streaming:

**Relaxation towards
the equilibrium**



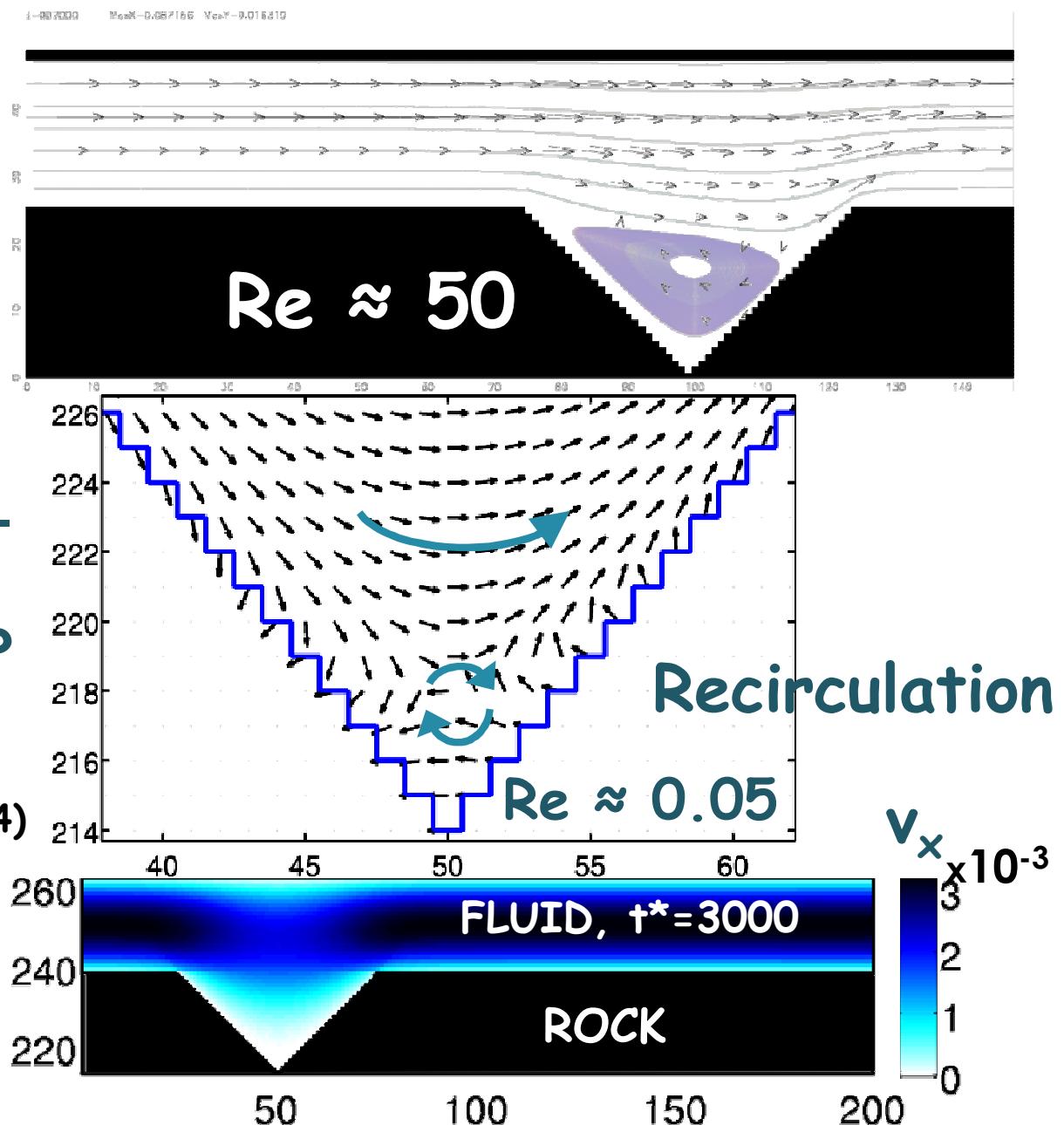
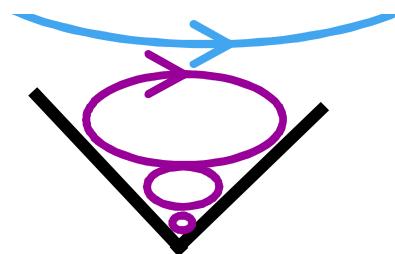
- Boltzmann methods:
 - > Average in a mesoscopic volume of particles occupation in a given direction
 - > Use of particles distributions : here, mass
 - > Conservation of mass and momentum

Flow Recirculation in a corner

Turbulence:
At large Re
Asymmetric

Moffat eddies:
At small Re
For angle < 146°
Symmetric

Moffat, J. Fluid Mech. (1964)



Hydro-thermal flows – meth. 1

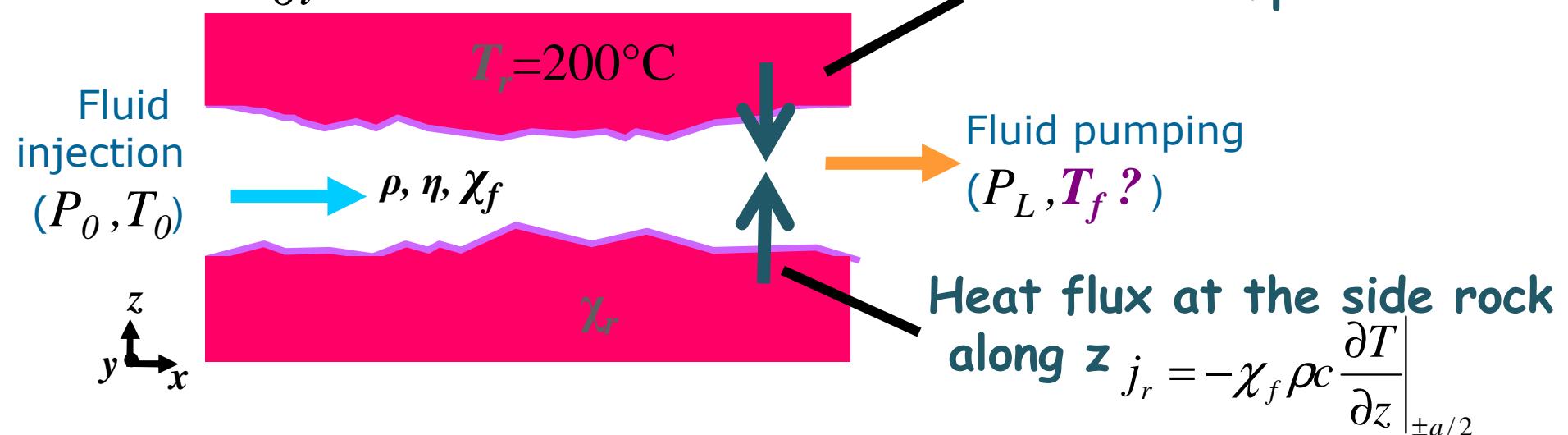
Finite differences

- Reynolds equation

$$\vec{\nabla} \cdot (a(x, y)^3 \vec{\nabla} p) = 0$$

- Heat diffusion-advection equation

$$\frac{\partial T}{\partial t} + \vec{\nabla} \cdot (\vec{v} T) = \vec{\nabla} \cdot (\chi \vec{\nabla} T)$$



Notations:

Fluid density: ρ

Dynamic viscosity: η

Fluid thermal diffusivity: χ_f

Rock thermal diffusivity: χ_r

Velocity: $\vec{v}(x, y, z)$

Averaged velocity: $\vec{u}(x, z)$

Pressure: $p(x, y, z)$

Hypotheses – Meth1

- Internal energy flux averaged across the fracture:

$$\vec{f} = \int_{a(x,y)} [E_0 + \rho c(T - T_0)] \vec{v} dz$$

$$\bar{T}(x, y) = \frac{\int_a^a v(x, y, z) T(x, y, z) dz}{\int_a^a v(x, y, z) dz}$$

- Hyp: $q_x \frac{\partial T}{\partial x} + q_y \frac{\partial T}{\partial y} = g(x, y)$

- Quartic profile: $T - T_r = -\frac{g}{32a^3 \chi_f} (4z^2 - a^2)(4z^2 - 5a^2)$

Thermal lubrication approximation

- Flux balance: $\vec{\nabla} \vec{f} + 2 j_r = 0$ with $j_r = -\chi_f \rho c \frac{\partial T}{\partial z} \Big|_{\pm a/2}$

- 2D temperature equation

$$\vec{q}(x, y) \cdot \vec{\nabla} \bar{T} + \frac{2\chi}{a(x, y)} Nu (\bar{T} - T_r) = 0$$

Notations (for the fluid):

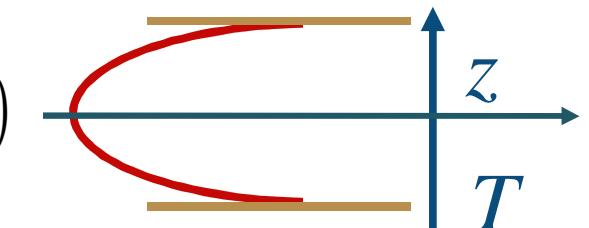
Density: ρ

Dynamic viscosity: η

Thermal diffusivity: χ_f

Specific heat capacity : c

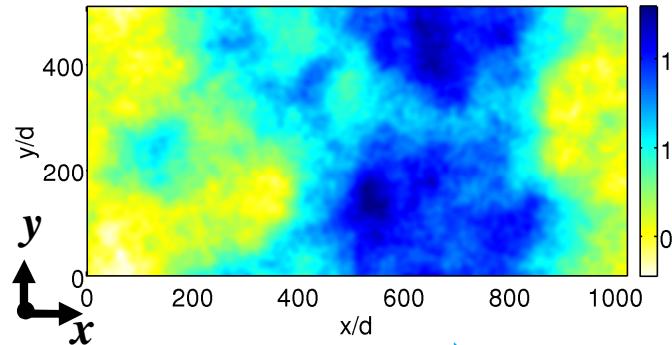
Velocity: $v(x, y, z)$



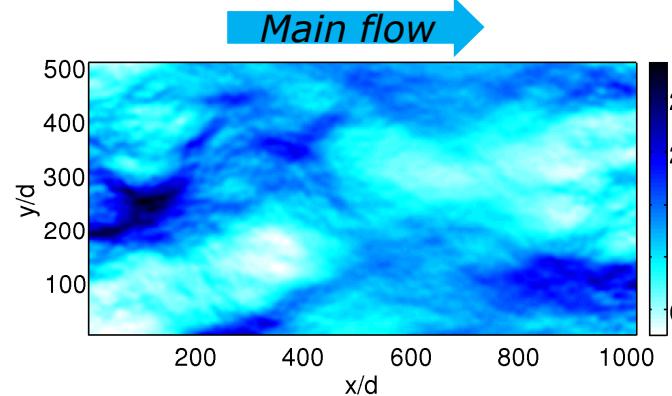
$$Nu = \frac{j_r}{j_{macro}} = \frac{70}{17}$$

Solving : Finite differences
+ biconjugate gradient method

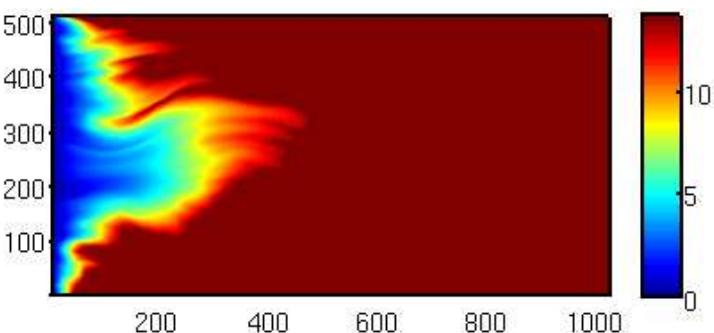
Illustration...hydro-thermal result



Rough apertures
 $a^*(x, y) = \frac{a(x, y)}{A}$



2D-flow norm
 $\|\vec{q}^*(x, y)\|$
 $= \frac{12\eta L_x \|\vec{q}(x, y)\|}{\Delta P A^3}$



$-\ln(\bar{T}^*)$
 $\bar{T}^* = \frac{\bar{T} - T_0}{\Delta T}$

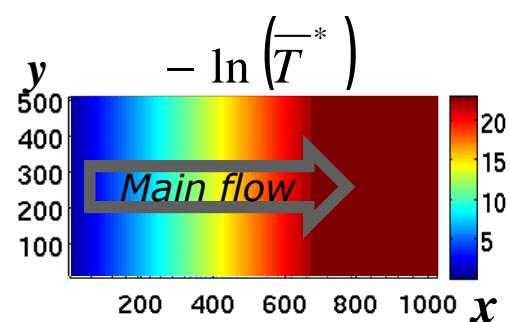
Example

A	1 mm
σ	0.35 mm
$L_x \times L_y$	1 x 0.5 m ²

Dyn. visc. η (10 bar; 100°C)	3.10 ⁻⁴ Pa.s
$\partial p / \partial x$	250 Pa/m
H	0.82 mm

Density ρ	1.10 ³ kg/m ³
ΔT	120° C
Fluid diffusivity χ	0.17 mm ² /s

Reference case: Parallel plates

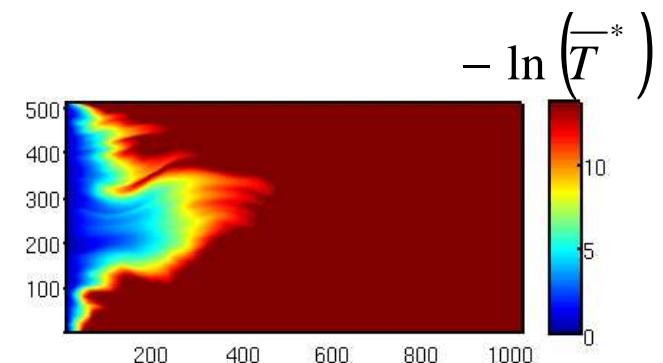
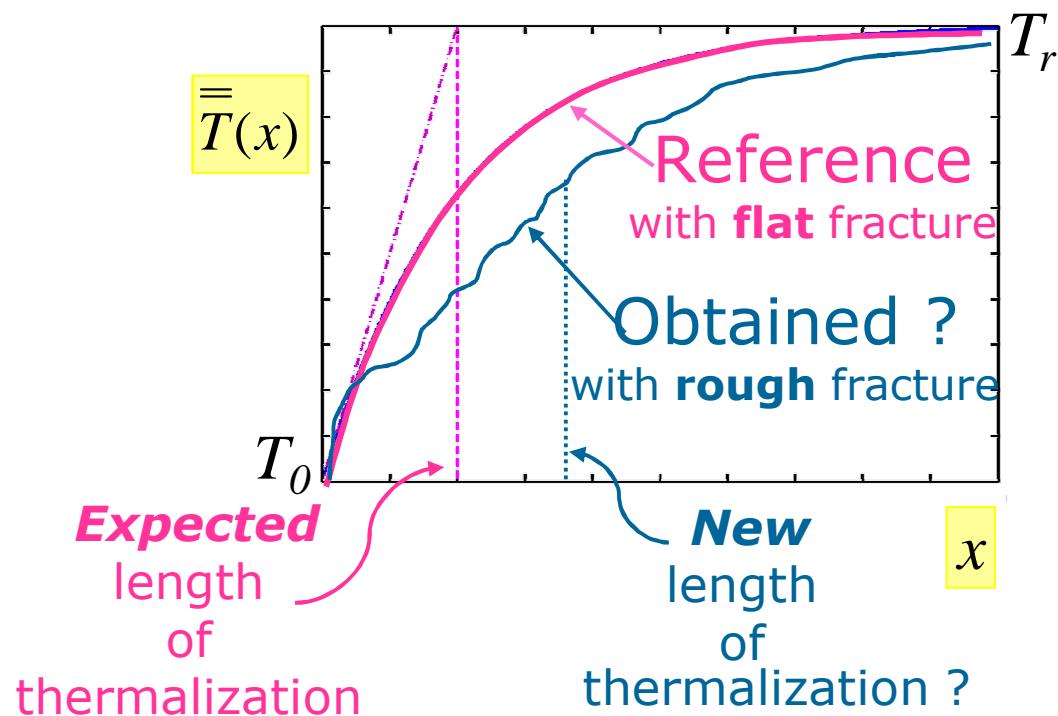


➤ *Analytic solution*

$$\bar{T}_{\parallel} - T_r = (T_0 - T_r) \exp\left(-\frac{x}{R_{\parallel}}\right)$$

$$R_{\parallel} = \frac{A^2 \|\vec{q}_{\parallel}\|}{2 Nu \chi_f}$$

○ Temperature in 1D

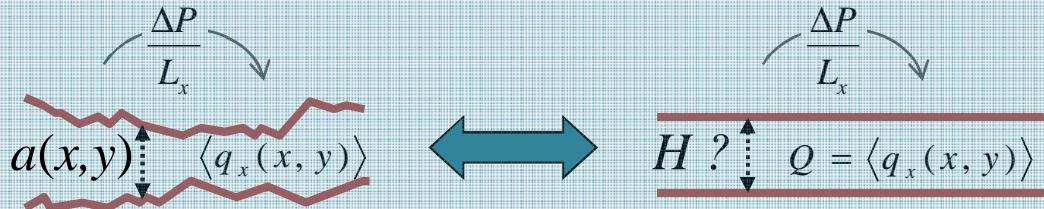


$$\bar{T}(x) = \frac{\int u_x(x, y) \bar{T}(x, y) dy}{\int u_x(x, y) dy}$$

$$u_x = \int_a v(x, y, z) dz / a(x, y)$$

Equivalent aperture definitions

> Hydraulic aperture H



$$\bar{q}(x, y) = \int_a^y \vec{v}(x, y, z) dy = -\frac{a^3}{12\eta} \vec{\nabla}_2 p$$

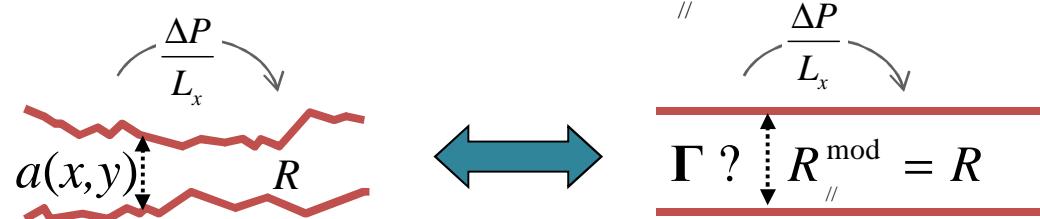
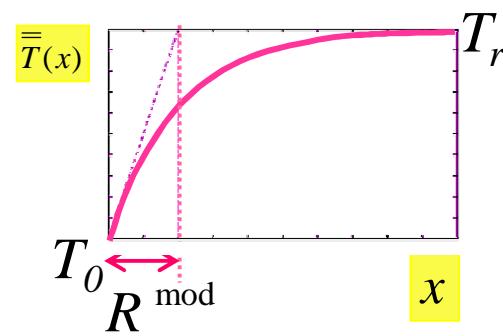
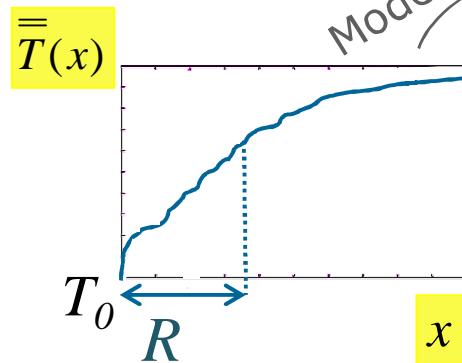
$$H^3 = -\langle q_x(x, y) \rangle 12\eta \frac{l_x}{\Delta P}$$

> Thermal aperture Γ

Parallel plates solution

Modeled with

$$\bar{T}_{//} - T_r = (T_0 - T_r) \exp\left(-\frac{x}{R_{//}^{\text{mod}}}\right)$$

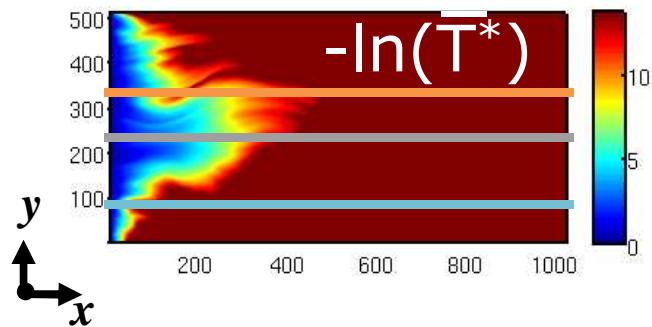


$R_{//}^{\text{mod}}$: Suitable thermal length

$$\Gamma \propto (R_{//}^{\text{mod}})^{1/4}$$

Illustration...thermal characterization

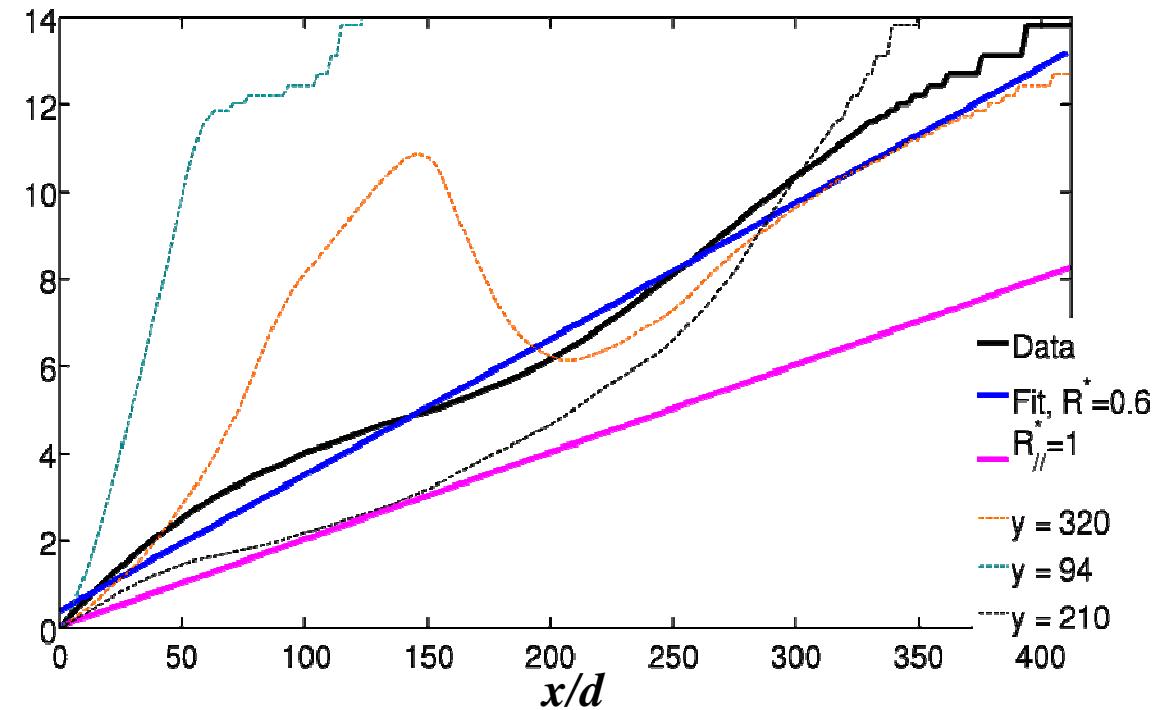
2D Temperatures



$\bar{\bar{T}}(x)$

$\bar{\bar{T}}_{\parallel}^{\text{mod}}$ with Γ
 $\bar{\bar{T}}_{\parallel}$ with A

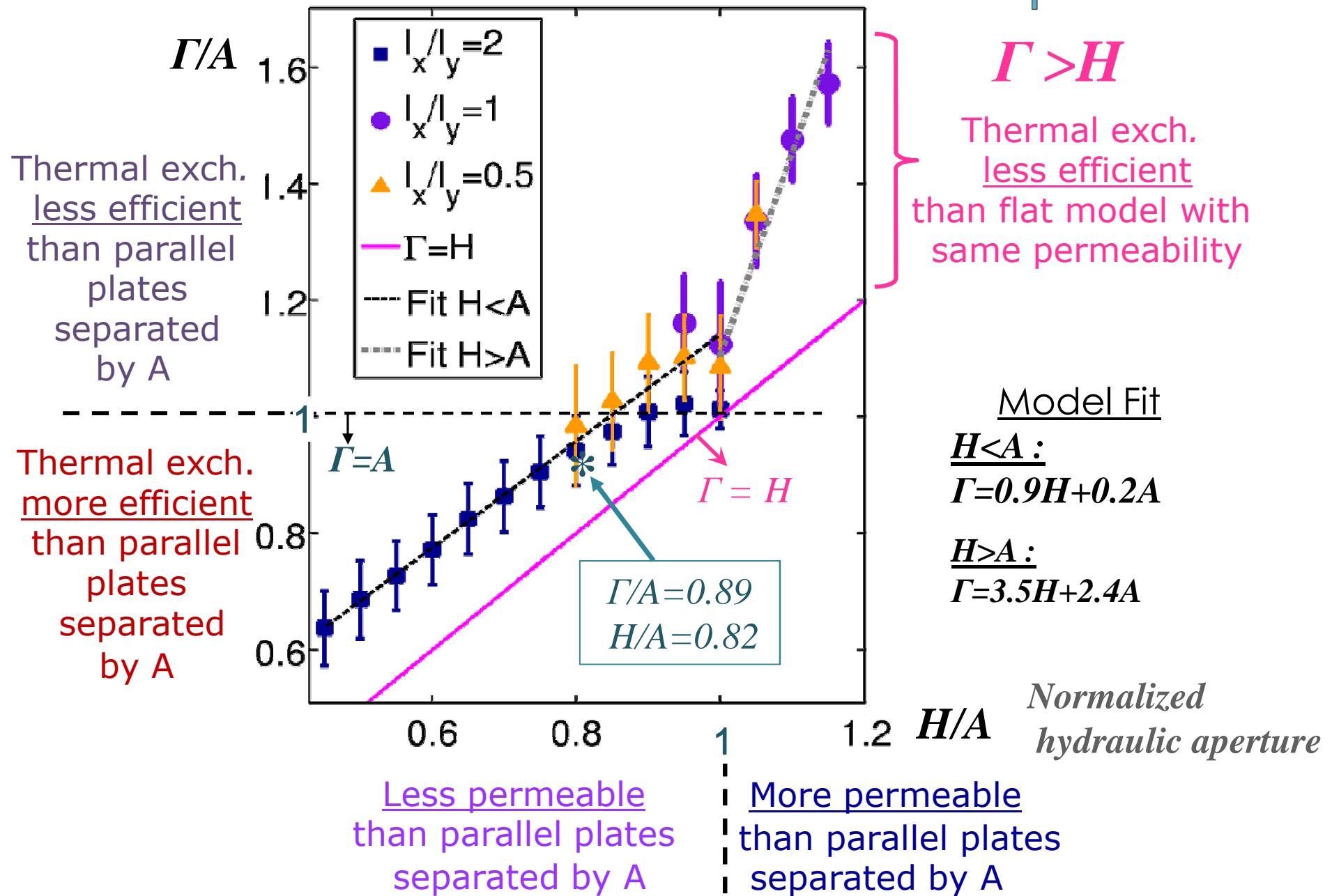
1D Temperatures $-\ln(\bar{\bar{T}}^*)$



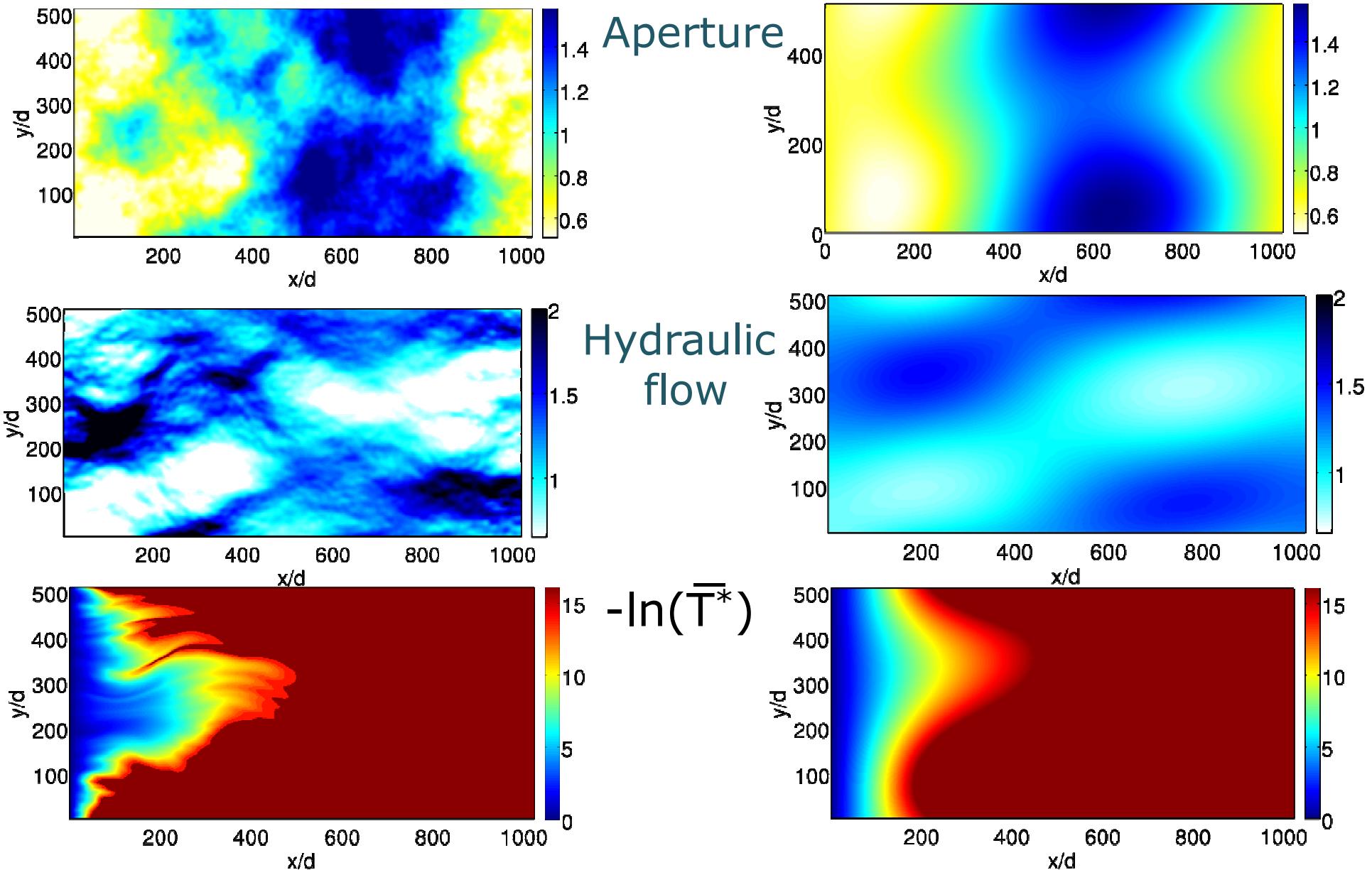
➤ Thermal aperture : $\Gamma = 0.89 \text{ mm} \dots < A = 1 \text{ mm}$

Enhanced
thermal behavior

Statistical results for thermal apertures



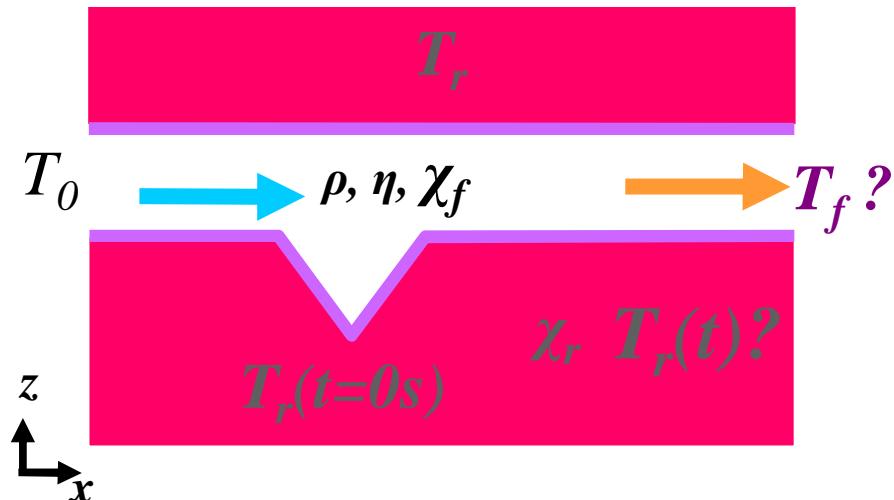
Control of the large scales modes on the hydro-thermal variations



Hydro-thermal equations (Meth 2)

- Off lubrication regime:

ρ : density; χ : Thermal diffusivity; η : Viscosity



Effect of a sharp morphology on

- > the fluid temperature ?
- > the rock temperature ?

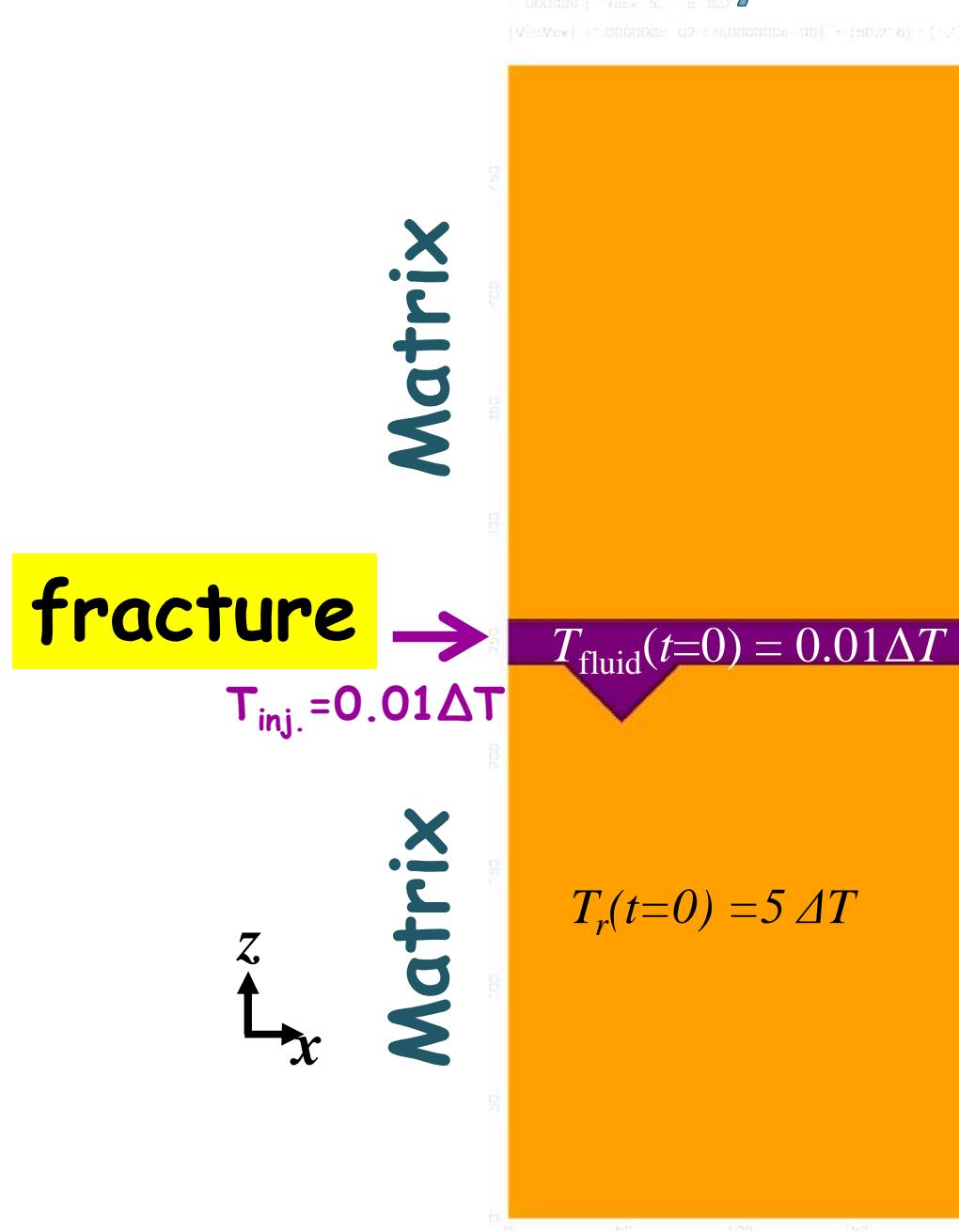
- Rock temperature variable in space and time

$$\frac{\partial T}{\partial t} + \vec{\nabla} \cdot (\vec{v} T) = \vec{\nabla} \cdot (\chi \vec{\nabla} T)$$

Solved in fluid and rock
(Thermal diffusivity χ_f and χ_r)

- Use of a second distribution particles with LB methods
- Conservation of internal energy and energy flux

Moderate Reynolds number



Space unit : $A/20$

Time unit Δt : $[A/(20)]^2 \cdot (1/\chi_r)$

Temperature unit: arbitrary (fixed by fluid injection temperature and rock temperature)

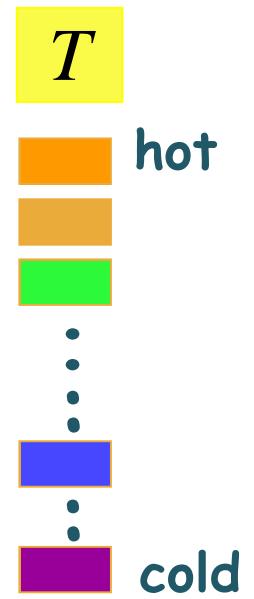
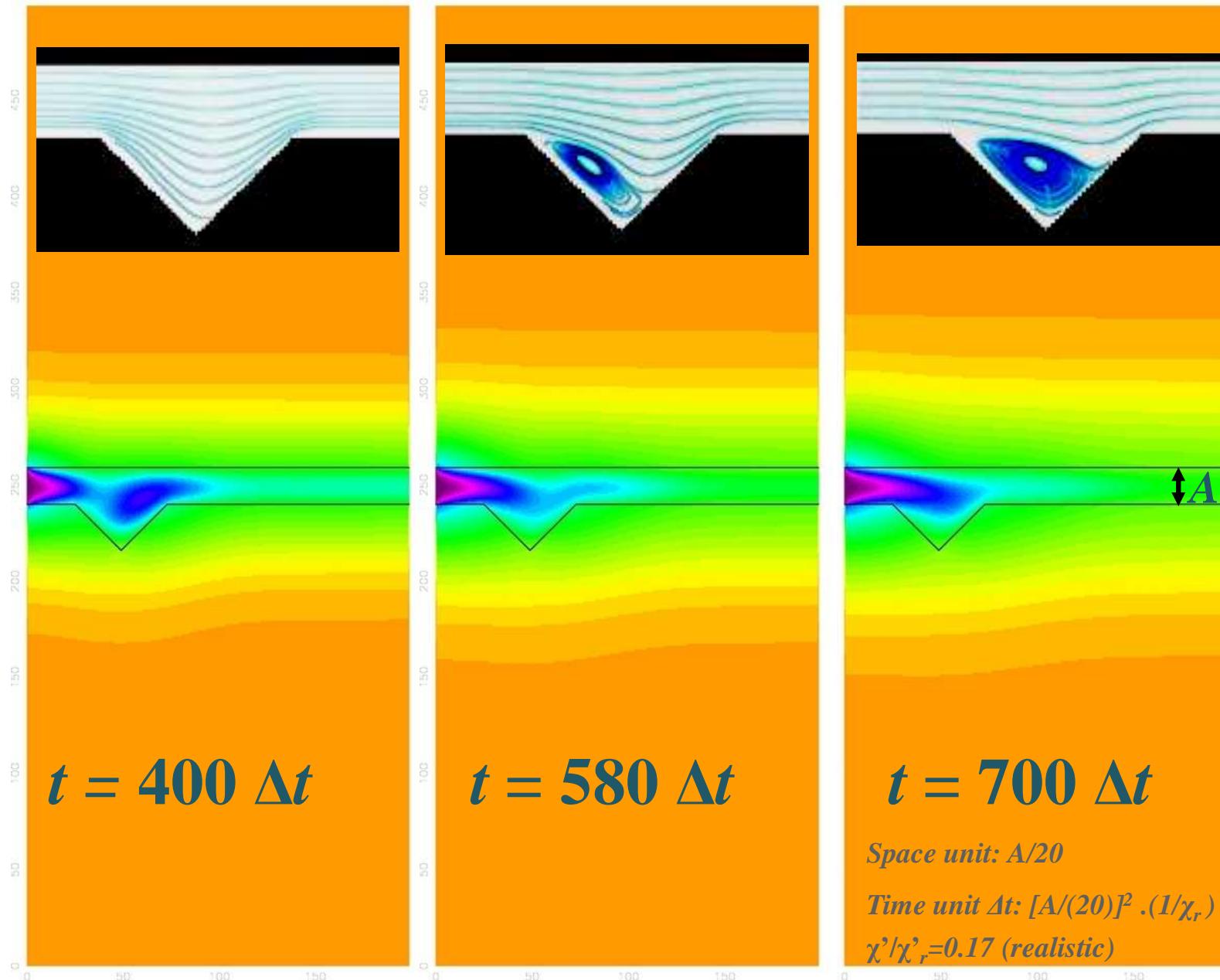
$\chi'/\chi'_r = 0.17$ (realistic)

$$\text{Peclet}=45 = \frac{\text{Convection}}{\text{Conduction}}$$

Colorscale:

Moderate Reynolds number

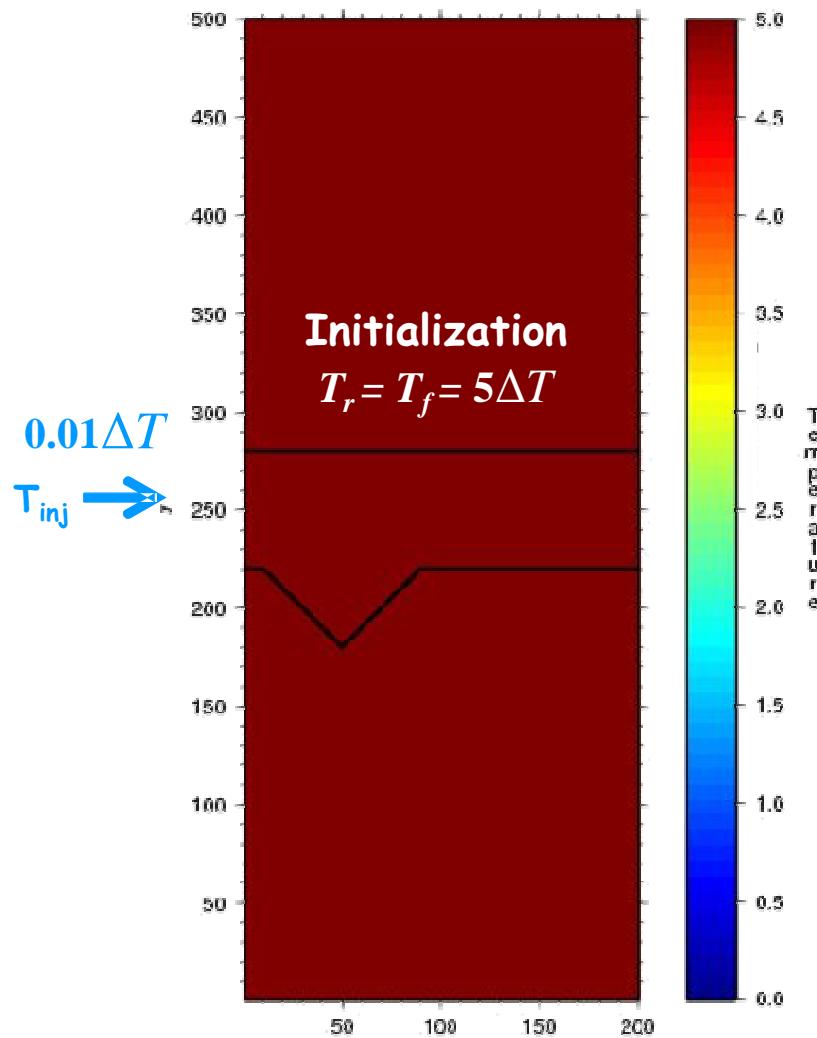
$t=0$; $\text{Re}=640$; $\text{Tr}=0.000000e+00$; $\text{Tu}=654$; $\text{Tr}_0=0.000000e+00$; $\text{Tu}_0=0.000000e+00$
 $[\text{Min}:\text{Max}] = [1.000000e-02 : 5.000000e+00] \text{ in } \langle 1,241 \rangle : \langle 1,1 \rangle$
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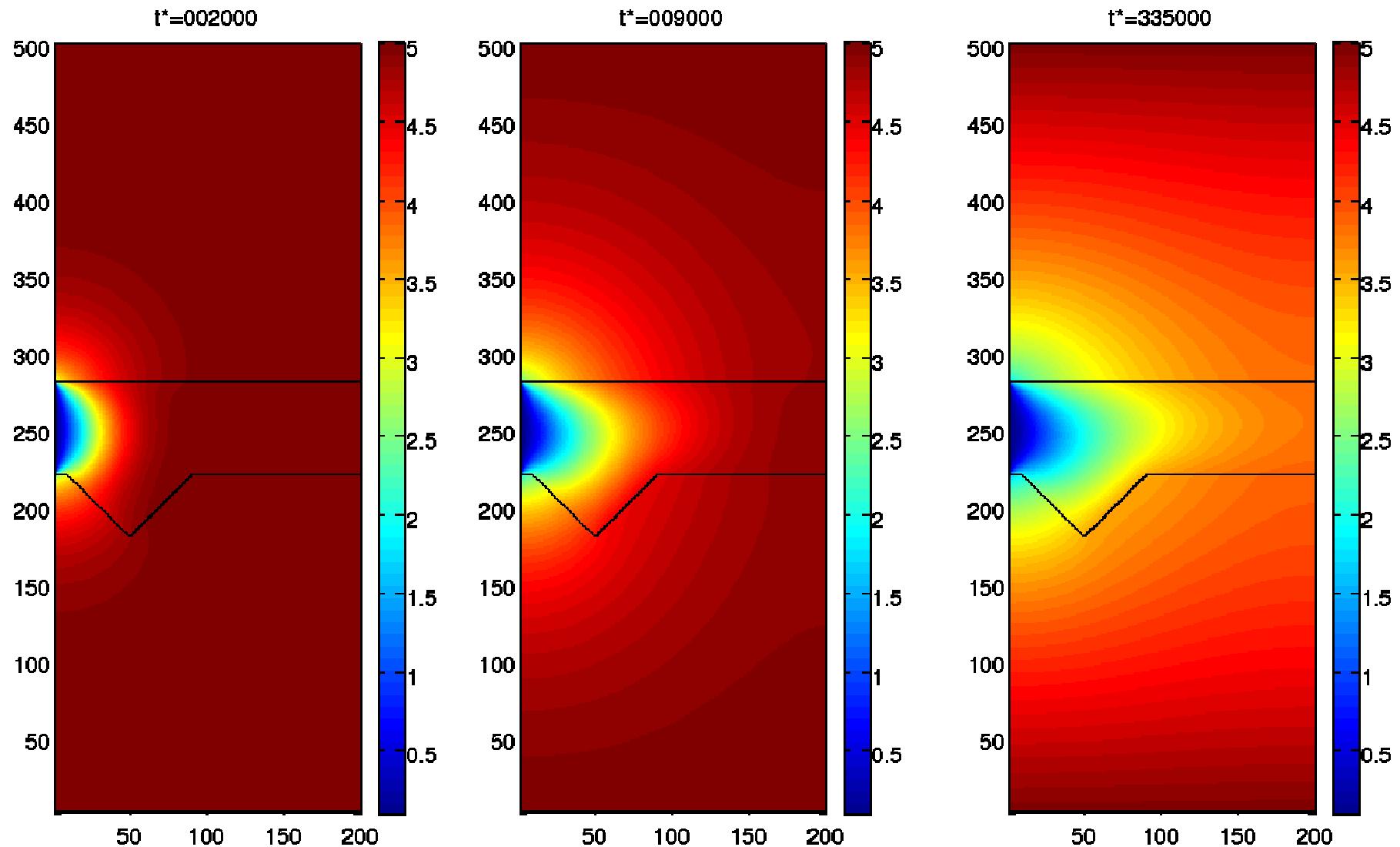
Large channel perturbed by a corner

• Initialization

- Channel of $200\Delta x \times 120\Delta x$ perturbed by a corner, inside rock
- Whole system size: $200\Delta x \times 500\Delta x$



Long term temperature field

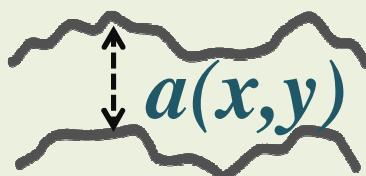


Conclusion and perspectives

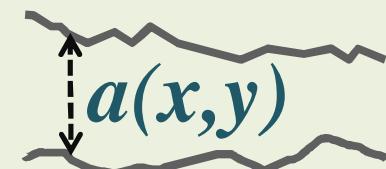
Draix cores and Draix permeability

- Methods developed to reconstruct fracture apertures from borehole core
- Characterization of
 - > surfaces
 - > Apertures
- Two models of aperture observed

Correlated
isotropic
surfaces



Indepedant
anisotropic
surfaces



- Permeability at core scale (10^{-9}m^2)
- Some larger scale hydraulic data are expected for comparison

Conclusion and perspectives

Hydro-thermal behavior under lubrication approx.:

- Due to roughness, channeling of
 - > Hydraulic flow
 - > Temperature (energy)
- Study of the aspect ratio L_x/L_y
- Large scale variations seems relevant
- Coarse grained behavior :
 - > Mechanical aperture A
 - > Hydraulic aperture H
 - > Thermal aperture Γ
- Thermal exchange less efficient than flat model with same permeability
- Laws proposed about H/A , Γ/A
- Integration in network modeling ?

Conclusion and perspectives

Hydro-thermal modeling with LB method

- Advantages
 - > Full hydraulic and heat equation solved in 3D
 - > Off lubrication regime
 - > Diffusion in the rock and liquid
- Long term behavior of geothermal systems
- How does
 - > Sharp morphology
 - > Moderate velocity

} change the thermal field ?
- Need to explore more parameters to draw a conclusion about the influence of an asperity with steep slopes !
- Characteristic length of scale of the recirculation ?
- Integrate more complex/realistic morphology for the rock